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# Electron spin resonance in CuCrS<sub>2</sub> chrome-copper disulphides synthesized by different methods

Galina Abramova,<sup>1,a)</sup> Anatolii Pankrats,<sup>1</sup> German Petrakovskii,<sup>1</sup> Julia C. E. Rasch,<sup>2</sup> Martin Boehm,<sup>2</sup> Aleksandr Vorotynov,<sup>1</sup> Vasilii Tugarinov,<sup>1</sup> Rita Szumszak,<sup>3</sup> Asya Bovina,<sup>1</sup> and Viktor Vasil'ev<sup>1</sup>

<sup>1</sup>L.V. Kirensky Institute of Physics, Russian Academy of Sciences, Siberian Branch, Akademgorodok 50, bld. 38, Krasnoyarsk 660036, Russia <sup>2</sup>Institute Max von Laue-Paul Langevin, Grenoble, Cedex 9, France

<sup>3</sup>Institute of Physics, Polish Academy of Sciences, 02-668 Warsaw, Poland

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The electron spin resonance (ESR) in CuCrS<sub>2</sub> disulphides is found to be strongly dependent on a synthesis method used. At a temperature of 300 K, a polycrystalline CuCrS<sub>2</sub> sample is paramagnetic with a g-value of 1.95 at 40 K, it undergoes the magnetic transition. In the temperature range 4.2–290 K, a single-crystal sample prepared by a chemical vapor transport method exhibits the ESR features typical of a ferromagnet. It is shown that these features are related to the presence of a small amount of the single-crystal CuCr<sub>2</sub>S<sub>4</sub> impurity in the CuCrS<sub>2</sub> single crystal. © 2010 American Institute of Physics. [doi:10.1063/1.3374679]

#### I. INTRODUCTION

The basic obstacle in the way of the development in thin-film multilayers for microelectronic applications is the necessity of using expensive technologies. The promising matrices for the creation of novel multilayers could be quasitwo-dimensional layered structures that represent threedimensional (3D) crystals with the strong anisotropy of chemical bonds. The use of intercalation and cation substitution make it possible to obtain low-cost magnetic/ nonmagnetic and dielectric/metal multilayers similar to heterostructures.

Layered transition-metal dichalcogenides and their intercalated complexes are of interest for practical applications, in particular, spintronic devices [see, for example (Ref. 1)] because of the features of their crystal structure<sup>2-4</sup> and physical properties.<sup>4-6</sup> These crystals are known to be highly anisotropic; most of them exhibit pronounced two-dimensionality resulting from the layered structure with weak bonds between the layers.<sup>4</sup> More precisely, the crystals consist of stacks of  $TX_2$  (T is the transition metal and X is the chalcogen) sandwiches separated by two layers of the empty tetrahedral positions called the van der Waals gaps, in which intercalant species (mostly, electron donors), such as alkali atoms, organic molecules, and 3D-transition-metal atoms can be inserted. A chrome-copper disulphide CuCrS<sub>2</sub> can be considered a Cu-intercalated dichalcogenide<sup>3,4</sup> with Cu ions intercalated into the van der Waals gaps between the CrS<sub>2</sub> layers.

At room temperature,  $CuCrS_2$  is a paramagnetic electron-ion semiconductor; at 670 K,<sup>7</sup> it undergoes the superionic conductor-semiconductor transition; at 110 K,<sup>8</sup> the electronic transition; and at 37 K,<sup>9,10</sup> the paramagnet-antiferromagnet transition. Neutron diffraction studies<sup>9,10</sup>

showed the incommensurate helimagnetic structure with the strong spin-phonon coupling below 37 K. The preliminary data we obtained previously<sup>5</sup> revealed negative magnetoresistance in some  $CuCrS_2$  solid solutions.

In this study, we present the results of the electron spin resonance (ESR) investigation of the  $CuCrS_2$  samples synthesized using different methods.

#### **II. EXPERIMENTAL**

The experiments were carried out on polycrystalline and single-crystal CuCrS<sub>2</sub> samples. Polycrystalline samples were prepared from pure Cu, Cr, and S elements by heating in vacuum quartz ampoules; CuCrS<sub>2</sub> single crystals were grown by a chemical vapor transport (CVT) method in vacuum quartz ampoules.<sup>9</sup> X-ray diffraction (XRD) spectra were recorded at room temperature with a DRON diffractometer with Cu  $K\alpha$  radiation. ESR spectra were obtained with a SE/X 2544 spectrometer at a frequency of 9.4 GHz and a computer-controlled magnetic resonance spectrometer with the operating frequency range 25–80 GHz in pulse magnetic field.

#### **III. RESULTS AND DISCUSSION**

The x-ray diffraction spectrum of the synthesized polycrystalline CuCrS<sub>2</sub> sample presented in (Fig. 1, curve 1) corresponds to the crystal structure with an R3m rhombohedral space group (the card 74–0196). Polycrystalline CuCrS<sub>2</sub> was found to be single-phase with the lattice parameters a= 3.480 Å and c ranged from 18.70 to 18.72 Å for different samples. The XRD spectrum (Fig. 1, curve 2) makes one conclude that the CVT crystal plate corresponds to the (001) plane of the CuCrS<sub>2</sub> R3m lattice with the c-axis perpendicular to the plate. In addition, the spectrum of the CVT CuCrS<sub>2</sub> single crystals contains an additional line corresponding to the CuCr<sub>2</sub>S<sub>4</sub> phase with a spinel lattice (Fd3m space group) and the lattice parameter a=9.79 Å.

<sup>&</sup>lt;sup>a)</sup>Author to whom correspondence should be addressed. Electronic mail: agm@iph.krasn.ru.





FIG. 1. XRD spectra of the CVT single-crystal and polycrystalline  $\rm CuCrS_2$  samples at 300 K.

The frequency-field dependences (FFDs) of magnetic resonance of the polycrystalline and single-crystal CVT CuCrS<sub>2</sub> samples taken at 90 K are presented in Fig. 2. The linear gapless FFD with a g-value of 1.95 (Fig. 2, curve 1) observable for the polycrystalline CuCrS<sub>2</sub> indicates the paramagnetic state of the sample at  $T > T_N$ . Similar paramagnettype magnetic resonance was found for blocking single crystals synthesized by crystallization from powder CuCrS<sub>2</sub> sulfide melt. However, the CVT CuCrS<sub>2</sub> single crystal was found to possess the strong anisotropy of the resonance properties. At a temperature of 90 K, different FFDs for parallel and perpendicular orientations of magnetic field relative to the CVT single crystal plate were obtained (see also Fig. 2). Angular dependences of the resonance field and linewidth measured at 300 K for the CVT CuCrS<sub>2</sub> single crystal are shown in Fig. 3. They indicate the anisotropy of the resonance properties of the sample at room temperature. The resonance fields at 300 K vary from 1.9 to 4.8 kOe upon



FIG. 2. FFDs of the magnetic resonance at 90 K for the polycrystalline and CVT single-crystal CuCrS<sub>2</sub> samples.

FIG. 3. Angular dependences of the resonance field (a) and linewidth (b) at 300 K for the CVT CuCrS<sub>2</sub> single crystal.

turning the magnetic field from the *ab*-plane to *c*-axis. The observed behavior of the resonance properties of the CVT CuCrS<sub>2</sub> single crystal is typical of an easy-axis ferromagnet. The difference between the FFDs of the polycrystalline and CVT single crystal samples can be explained by the presence of the ferromagnetic ordered state with easy and hard magnetization directions parallel and perpendicular to the CVT single crystal plate, respectively. Obviously, the ferromagnetic properties of the CVT CuCrS<sub>2</sub> crystal are connected with an additional CuCr<sub>2</sub>S<sub>4</sub> impurity phase (5%–20% for different samples) discovered by an XRD analysis and confirmed by the temperature dependences of the magnetization and resonance fields of the CVT CuCrS<sub>2</sub> single crystal (Fig. 4).

Figure 4(a) depicts temperature dependences of the magnetization for the polycrystalline and CVT CuCrS<sub>2</sub> single crystal. The temperature of the magnetic transition in CuCrS<sub>2</sub> depends on a synthesis method used and varies within 37-40 K. A value of CuCrS<sub>2</sub> magnetic susceptibility also changes for different samples but the linear field dependence of the magnetization measured at 300 K for the polycrystalline sample in applied magnetic fields up to 50 kOe [inset in Fig. 4(a) indicates the paramagnetic state of the CuCrS<sub>2</sub> sample at room temperature. One can see in Fig. 4(a) that the magnetization of the CVT single-crystal CuCrS<sub>2</sub> (curve 2) contains an additional magnetic contribution (curve 3), which can be related to the presence of ferromagnetic CuCr<sub>2</sub>S<sub>4</sub> with the Curie temperature close to 400 K. Temperature dependences of the resonance field taken at 35 GHz for parallel and perpendicular orientations of magnetic field relative to the CVT single-crystal  $CuCrS_2$  plate are given in Fig. 4(b). The difference between the resonance fields for the two orientations of the external magnetic field proportional to the magnetic anisotropy of the impurity phase decreases with the increasing temperature and vanishes at approaching a temperature of about 400 K, which is consistent with the Curie temperature of  $CuCr_2S_4$  (Ref. 11).



FIG. 4. Temperature dependence of magnetic susceptibility for the polycrystalline (1) and CVT single-crystal (2) CuCrS<sub>2</sub> samples (a) measured at the different magnetic fields. Inset: field dependence of magnetization for the polycrystalline sample at 300 K. Temperature dependences of the magnetic resonance field for the external magnetic field directed along *c*-axis and in *ab*-plane of the CVT single crystal (b).

It should be emphasized that both the host  $CuCrS_2$  and the impurity phases in the CVT  $CuCrS_2$  samples are singlecrystal. The presence of the single-crystal  $CuCr_2S_4$  impurity in the rhombohedral  $CuCrS_2$  single crystal may be attributed to the similarity of their crystal structures. The structural spinel-rhombohedral transition is characteristic of some Cucontaining compounds, such as  $CuTi_2S_4$  (Ref. 12) and  $CuFeO_2$ .<sup>13</sup> The description of this transition and the comparison of the spinel and rhombohedral lattices for  $CuFeO_2$  were presented in Ref. 13.

The  $Cu^{2+}$  and  $Cr^{3+}$  ions in the spinel  $CuCr_2S_4$  lattice occupy tetra-sites and octa-sites, respectively, similar to the R3m lattice of CuCrS<sub>2</sub>. Both the spinel and rhombohedral lattices can be presented as a layer sequences of the tetrahedral and octahedral sites along the [001] and [111] directions for the R3m and Fd3m structures, respectively (Fig. 5). Both the CuCrS<sub>2</sub> and CuCr<sub>2</sub>S<sub>4</sub> lattices contain similar layers of the octahedral sites occupied by Cr<sup>3+</sup>-ions. The length of chemical bonds between Cr and S ions in  $CuCr_2S_4$  is close to 2.34 Å. The Cr-S chemical bond length in CuCrS<sub>2</sub> lies within 2.42-2.38 Å and strongly depends on a synthesis method used. We suggest the Cr<sup>3+</sup> layers in the CVT crystal marked by arrow in the picture to be common for two adjacent layer stacks with the different crystal structures. Thus, the CuCrS<sub>2</sub> solid solutions can be perspective matrices for synthesis of single-crystal heterostructures.

#### **IV. CONCLUSIONS**

The  $CuCrS_2$  single crystals prepared by the CVT possess the intriguing resonance properties, which may be attributed



FIG. 5. (Color online) Comparison of the  $CuCrS_2$  (s.g. R3m) and  $CuCr_2S_4$  (s.g. Fd3m) crystal lattices.

to a small amount of the  $\text{CuCr}_2\text{S}_4$  impurity phase. In the temperature range 90–300 K, the ESR spectra of the samples are typical of ferromagnets and the FFDs vary with the external magnetic field orientation. The obtained temperature and FFDs of the ESR suggest that the features of the resonant properties of the crystal are related to the  $\text{CuCr}_2\text{S}_4$  impurity layers with the single-crystal structure. As the XRD data show, the  $\text{CuCr}_2\text{S}_4$  [111] crystal planes are parallel to the  $\text{CuCrS}_2$  [001] planes. In the temperature range 90–300 K, the resonance and magnetic properties of pure  $\text{CuCrS}_2$  are typical of paramagnets.

The presence of the  $CuCr_2S_4$  impurity in the chromecopper disulphide is caused by the intercalation ability of the  $CuCrS_2$  crystal structure. This property can be used to prepare heterostructures with single-crystal layers of  $CuCr_2S_4$  or another magnetic impurity between the single-crystal  $CuCrS_2$  layers.

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