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# Synthesis, structural and vibrational properties of microcrystalline $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub>

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## ABSTRACT

Low-temperature rubidium samarium dimolybdate,  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub>, microplates have been prepared by solid state synthesis at 723 K by 70 h and 893 K by 150 h. The crystal structure of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> has been first refined by the Rietveld method in the space group *Pbcn* with unit cell parameters  $a=5.1431(2)$  Å,  $b=18.8195(7)$  Å, and  $c=8.1641(3)$  Å ( $R_B=4.82\%$ ). The crystal structure of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> consists of the layers of MoO<sub>4</sub> tetrahedrons corner-sharing with SmO<sub>8</sub> square antiprisms. About 20 narrow Raman lines have been observed in Raman spectra recorded for the  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> powder sample.

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## 1. Introduction

Molybdate crystals have become of considerable interest because of their interesting structural, luminescent and spectroscopic properties, promising wide applications in optical technologies [1–8]. Among the molybdates reported, the rare-earth containing crystals are of special attention because of their possible application as laser host materials [3,7,9–14]. In many such crystals the rare-earth ions are found in low symmetry positions, which is a key factor for creation of effective ultrafast laser gain media. Searching for new optical materials, it is reasonable to explore the long-discovered but less studied complex molybdate crystals because only fragmentary information can be found and potentials of the crystals are unclear in many aspects [15].

The low-temperature orthorhombic modification  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> is found in the quasi-binary system Rb<sub>2</sub>MoO<sub>4</sub>–Sm<sub>2</sub>(MoO<sub>4</sub>)<sub>3</sub> at Rb:Sm=1:1, and it exists below  $T=1183$ – $1203$  K [16,17]. At  $T=1183$ – $1203$  K,  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> transforms to a high-temperature modification  $\alpha$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> of an unknown symmetry. The crystal structure and physical parameters of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> remains unknown. The present study is aimed at the synthesis and evaluation of microstructural and vibrational parameters of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub>. The two-stage solid state synthesis at comparatively low temperatures is

selected for  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> preparation to provide precise chemical composition control.

## 2. Experimental

The powder samples of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> were prepared using the solid state reaction in platinum crucible. Analytically pure MoO<sub>3</sub> (99.9%), Rb<sub>2</sub>CO<sub>3</sub> (99.99%), and Sm<sub>2</sub>O<sub>3</sub> (> 99.9%) were used as starting materials. To remove the residual water occasionally captured from the air, the Rb<sub>2</sub>CO<sub>3</sub> was previously annealed at  $T=623$  K. Initially, rubidium and samarium molybdates were prepared. The heat treatment of stoichiometric mixtures was started at  $T=723$  K and followed by a step-wise temperature increase up to  $T=873$  K (Rb<sub>2</sub>MoO<sub>4</sub>) and 1073 K (Sm<sub>2</sub>(MoO<sub>4</sub>)<sub>3</sub>), respectively. After cooling to room temperature, the Rb<sub>2</sub>MoO<sub>4</sub> and Sm<sub>2</sub>(MoO<sub>4</sub>)<sub>3</sub> products were grinded and mixed to obtain the RbSm(MoO<sub>4</sub>)<sub>2</sub> composition. Then, the powder mixture of the components was preheated at 723 K for ~70 h and then it was fired at 893 K for 150 h to yield the final product. The micromorphology of the powder was observed by SEM with an LEO 1430 device. The measurements, however, were possible only at low accumulation times because of drastic surface charging of the RbSm(MoO<sub>4</sub>)<sub>2</sub> particles.

The diffraction data for Rietveld analysis were collected over the range of  $2\theta$ : 5°–100° at 298 K with a Bruker D8 ADVANCE powder diffractometer in the Bragg–Brentano geometry and linear Vantec detector. The operating parameters were: CuK $\alpha$  radiation, step size

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0.016°, and counting time 1.2 s per step. The X-ray patterns of the title compound were indexed using the ITO program [18]. Almost all reflections were indexed in orthorhombic space group *Pbcn* with cell parameters  $a=5.145$  Å,  $b=18.823$  Å, and  $c=8.167$  Å (FOM=45), excepting several minor intensity peaks that could not be assigned definitely. The structure of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> is isostructural to TlPr(MoO<sub>4</sub>)<sub>2</sub> [19] and we use the atomic coordinates to refine the structure of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub>. All refinements and data processing have been performed by DDM program [20]. The parameters of refinement are reported in Table 1S.

Unpolarized Raman scattering spectra were recorded from the powder sample using a triple grating spectrometer TriVista 777 and a line of  $\lambda=532$  nm (200 mW) at room temperature. The spectral resolution during measurements was  $\sim 1$  cm<sup>-1</sup> (FWHM). For comparison, a sample of  $\beta$ -RbNd(MoO<sub>4</sub>)<sub>2</sub> was recorded together with  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub>. The conditions of  $\beta$ -RbNd(MoO<sub>4</sub>)<sub>2</sub> preparation can be found elsewhere [10].

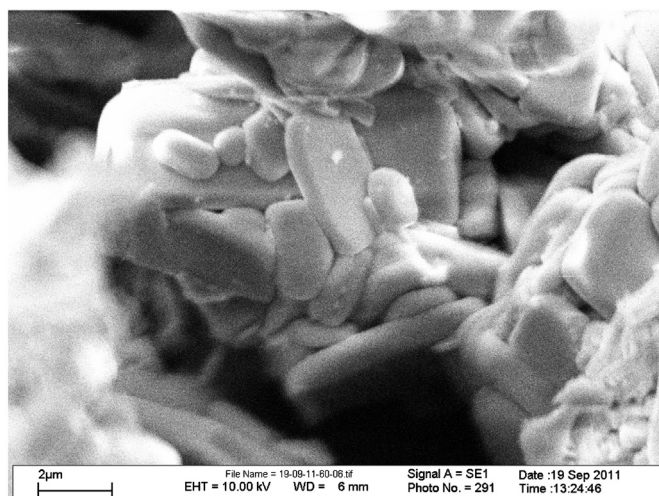


Fig. 1. SEM image of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> microplates.

### 3. Results and discussion

The final powder product was of white color with the light orange tint typical of samarium oxides. The micromorphology of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> particles is shown in Fig. 1. The powder is formed by uniform roundish coalescent grains of  $\sim 1$ – $4$   $\mu$ m in diameter. The big facets of the grains are very flat and this indicates the initial stage of equilibrium crystal shape formation. The refinement of the structure in space group *Pbcn* led to a minimal *R*-factor. The coordinates of atoms, isotropic thermal parameters and occupations of atom positions are shown in Table 2S. The experimental (dots) and theoretical (lines) X-ray diffraction patterns of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> are shown in Fig. 2. Evidently, a good relation between experimental and theoretical curves is achieved. The crystal structure of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> is shown in Fig. 3 [21]. The structure is formed by complex layers of MoO<sub>4</sub> tetrahedrons shared with SmO<sub>8</sub> square antiprisms by corners. The layers are orthogonal to the *b*-axis of the unit cell and this is a strongly selected crystallographic direction. So, the big flat surfaces of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> microplates visible in Fig. 1 seem to be orthogonal to the *b* crystallographic axis. Rubidium atoms are located in the gap between MoO<sub>4</sub>-SmO<sub>8</sub> layers and are coordinated by six O ions. A set of main interatom bond lengths is presented in Table 3S. In the  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> structure each Mo ion has four Mo–O bonds with comparatively short lengths:  $L(\text{Mo–O})=1.72$ – $1.85$  Å, and two very long Mo–O bonds with as high lengths as  $L(\text{Mo–O})=2.77$ – $2.79$  Å. So, the MoO<sub>6</sub> unit may be considered as a highly distorted octahedron.

Raman spectra of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> and  $\beta$ -RbNd(MoO<sub>4</sub>)<sub>2</sub> crystals are shown in Fig. 4. Raman spectra of both crystals are similar and consist of about 20 narrow Raman lines over the spectral range of 30–1000 cm<sup>-1</sup> (Table 4S). Thus, in the case of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub>, in the range of stretched vibrations of MoO<sub>n</sub> polyhedra (800–1000 cm<sup>-1</sup>), the five lines observed have a higher Raman shift. The intensive lines in this spectral range were also found in other complex molybdates [4,7,22–28]. In the stretched Mo–O vibrations range the difference between Raman line positions varies from 0.2% to 0.6% for  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> and  $\beta$ -RbNd(MoO<sub>4</sub>)<sub>2</sub>. Also, the Raman lines of stretched Mo–O vibrations are narrower in  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub>. Moreover, in the case of RbNd(MoO<sub>4</sub>)<sub>2</sub>, the spectral shape of 943 cm<sup>-1</sup> (the most intensive stretched Mo–O) line is well

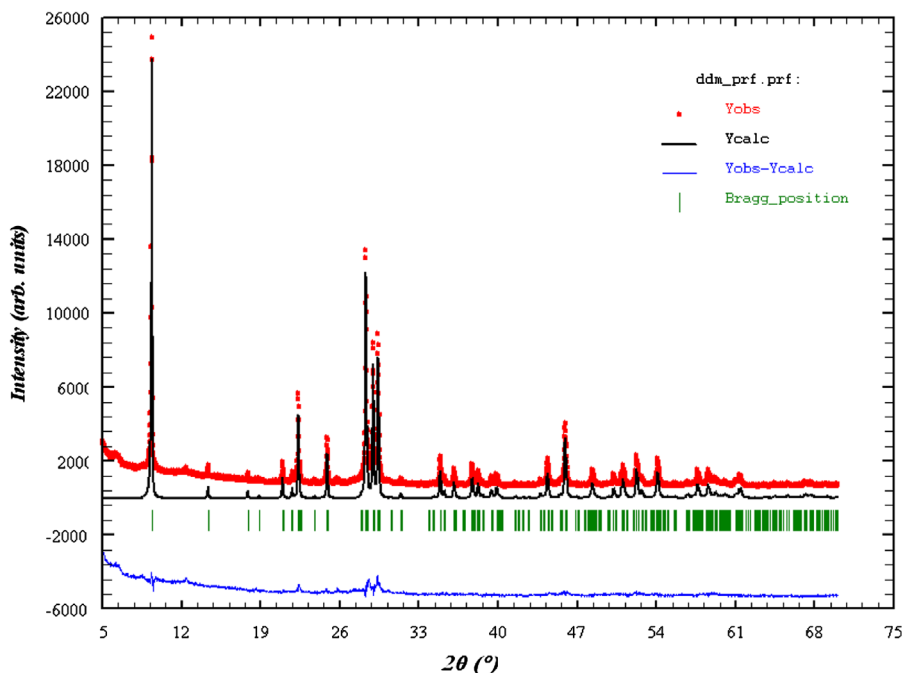


Fig. 2. XRD pattern recorded from the powder sample of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> at room temperature.

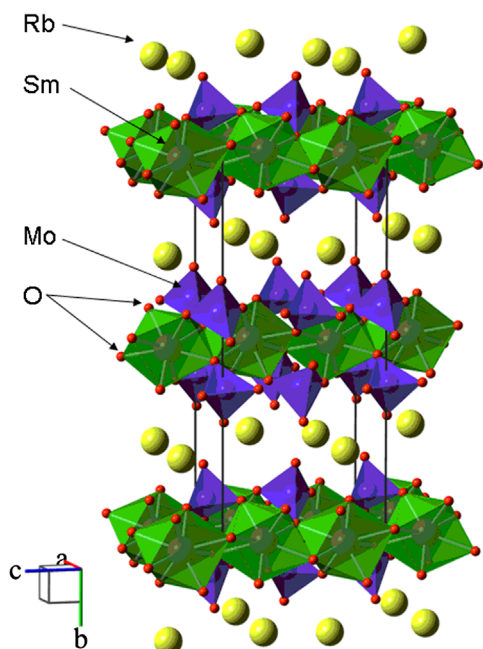


Fig. 3. Crystal structure of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub>. Unit cell is outlined. Lone samarium, molybdenum and oxygen atoms are omitted for clarity.

microcrystal growth. The layered crystal structure refined by the Rietveld method well relates to this supposition. The low symmetry positions are found for Sm<sup>3+</sup> ions. The spectroscopic parameters of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> verify the tetrahedral coordination of Mo<sup>6+</sup> ions and layered character of the crystal lattice. It should be pointed that, according to the phase diagram of the Rb<sub>2</sub>MoO<sub>4</sub>–Sm<sub>2</sub>(MoO<sub>4</sub>)<sub>3</sub> system, the crystal growth of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> is possible below  $T \sim 1200$  K [16,17] to avoid the  $\alpha \leftrightarrow \beta$  phase transition. In this situation, the top seeded solution growth (TSSG) method seems to be optimal.

### Acknowledgments

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### Appendix A. Supporting information

Supplementary data associated with this article can be found in the online version at <http://dx.doi.org/10.1016/j.matlet.2013.04.039>.

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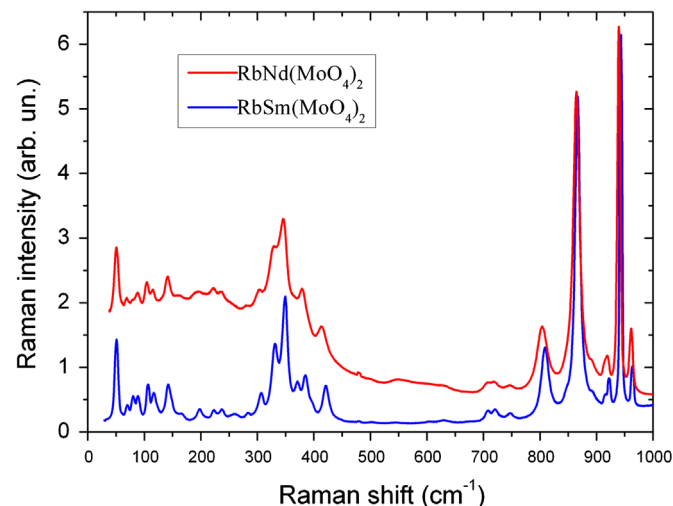


Fig. 4. Raman spectra of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> and  $\beta$ -RbNd(MoO<sub>4</sub>)<sub>2</sub>.

described by a Voigt contour with Gaussian width of 5 cm<sup>-1</sup> and Lorentzian width of 2 cm<sup>-1</sup>, while in the case of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> the spectral shape of 943 cm<sup>-1</sup> line is rectangular-like with width of 5 cm<sup>-1</sup>. Thus, the higher vibrational frequency and narrower lines in the case of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> allow us to conclude that Mo–O vibrations are more harmonic in this crystal in comparison with that of  $\beta$ -RbNd(MoO<sub>4</sub>)<sub>2</sub>. In the low-frequency part of Raman spectrum, the same tendency is observed and the Raman lines recorded from  $\beta$ -RbNd(MoO<sub>4</sub>)<sub>2</sub> crystal have a lower Raman shift in comparison with that of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub>. The observed difference between the counterpart lines is about 0.5–1 cm<sup>-1</sup>.

### 4. Conclusions

In the present study, binary molybdate  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> was prepared by solid state synthesis. The plate-like microcrystals of  $\beta$ -RbSm(MoO<sub>4</sub>)<sub>2</sub> seems to be generated due to anisotropic

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