



Thermite synthesis and characterization of Co–ZrO₂ ferromagnetic nanocomposite thin films



V.G. Myagkov^{a,*}, V.S. Zhigalov^{a,b}, L.E. Bykova^a, S.M. Zharkov^{a,c}, A.A. Matsynin^{a,c}, M.N. Volochaev^{a,b}, I.A. Tambasov^a, G.N. Bondarenko^d

^a Kirensky Institute of Physics, Russian Academy of Sciences, Siberian Branch, Krasnoyarsk, 660036, Russia

^b Reshetnev Siberian State Aerospace University, Krasnoyarsk, 660014, Russia

^c Siberian Federal University, Krasnoyarsk, 660041, Russia

^d Institute of Chemistry and Chemical Technology, Russian Academy of Sciences, Siberian Branch, Krasnoyarsk, 660049, Russia

ARTICLE INFO

Article history:

Received 28 September 2015

Received in revised form

21 December 2015

Accepted 30 December 2015

Available online 5 January 2016

Keywords:

Thermite reactions

Reactive films

Ferromagnetic nanocomposite films

ZrO₂

ABSTRACT

Co–ZrO₂ ferromagnetic nanocomposite thin films were successfully synthesized using a new thermite reaction between Zr and Co₃O₄ in layer geometry. The initial Zr/Co₃O₄ bilayers were obtained by the deposition of Zr layers onto Co₃O₄ films at room temperature. The process of mixing at the Zr/Co₃O₄ interface and synthesis of fine-crystalline Co and amorphous ZrO₂ phases started at a temperature above the initiation temperature $T_{in} \sim 250$ °C which did not depend on the bilayer thickness. For the bilayer thickness more than 300 nm high-temperature fcc-Co and cubic c-ZrO₂ (or tetragonal t-ZrO₂) phases were formed. For the bilayer thickness less than ~50 nm stable low-temperature hexagonal hcp-Co and monoclinic m-ZrO₂ phases were also present in the reaction product. A partial transformation from Co₃O₄ to CoO occurred after annealing at 300 °C. The secondary reaction between CoO and Zr started soon after Co₃O₄ had been converted to the CoO phase. After annealing at 500 °C more than 80% of Co was reduced and the final product contained Co nanoparticles above and below the superparamagnetic critical size embedded into a dielectric ZrO₂ matrix. The synthesized Co–ZrO₂ nanocomposite films possessed soft magnetic behavior, high magnetization and good chemical stability.

© 2016 Elsevier B.V. All rights reserved.

1. Introduction

Granular solids composed of ferromagnetic nanoparticles embedded in semiconducting (e.g. In₂O₃, TiO₂, ZnO, SnO₂) and insulating (e.g. SiO₂, Al₂O₃, MgO, HfO₂, ZrO₂) matrices have attracted a great deal of attention, both from the point of view of the fundamental perspective and for their potential applications [1–3]. These nanocomposites display magnetic and physicochemical properties which strongly depend on the method of preparation, the particle size, concentration and chemical bonding between the nanoparticles and matrix. Many approaches have been developed to synthesize ferromagnetic nanoparticles in semiconducting and insulating matrices such as the sol–gel method [4–6], rf magnetron sputtering [7–9], thermal decomposition [10], ball milling [11,12], etc. Among the above mentioned oxides ZrO₂ is a well known ceramic material which possesses excellent thermal,

electric, mechanical, optical properties, good chemical stability and has many practical applications [13–17]. The fabrication of ZrO₂-based ferromagnetic nanocomposites is a challenging goal, as this can provide ZrO₂ with the magnetic properties, which considerably enhance possibilities for applications. In the literature a major part of the papers is directed towards elucidating the nature of magnetism in transition metal doped ZrO₂ [18–20] and stabilization of the high temperature crystalline tetragonal (t-ZrO₂) and cubic (c-ZrO₂) phases [21–23]. Although ZrO₂-based magnetic nanocomposites obtained by various methods have been widely investigated, only a few reports are directed towards investigating magnetic, optical, magneto-optical, AC conductance and magnetoresistance properties in Co–ZrO₂ granular thin films [24–28]. The only method currently used for the fabrication of Co–ZrO₂ granular thin films is pulsed laser deposition. These films had an average Co nanoparticle size of about 2.5–63 nm against the Co volume fraction and showed a very sharp interface between the crystalline Co nanoparticles and the amorphous ZrO₂ matrix, with no evidence of intermixing [24–26]. Increasing the annealing temperature up to

* Corresponding author.

E-mail address: miagkov@iph.krasn.ru (V.G. Myagkov).

400 °C lead to the modification of the granular structure, which resulted in both the nucleation of ultra-small particles and the coalescence of neighboring particles [27]. In addition, J. Shin et al. reported the formation of magnetic nanostructures from self-assembled Co nanodots and nanowires embedded in a yttria-stabilized zirconia matrix which had been grown on Si(001) at 800 °C and 5×10^{-4} Torr O₂ pressure [28].

Recently, we have demonstrated a new thermite method for synthesizing ferromagnetic Fe–In₂O₃ [29] and Fe–ZrO₂ [30] nanocomposite films containing α -Fe clusters embedded into a transparent conducting In₂O₃ oxide and dielectric ZrO₂ matrix, respectively. In this report, we extend this approach to the formation of magnetic Co–ZrO₂ nanocomposites using a new thermite reaction (1):



Thermite reactions between fuels such as Al, Be, Mg, B, Hf, Li, Ta, Ti, Zr and oxidizers such as SiO₂, SnO₂, WO₃, MoO₃, CuO, Bi₂O₃, I₂O₅, Cr₂O₃, Fe₂O₃ react with a release of large quantities of heat, may have a self-propagating high-temperature synthesis (SHS) mode and widely used for the synthesis of materials. For traditional thermites which contain micron-sized particles, the combustion velocity is of the order of several centimeters to a few meters per second [31]. When a fuel and oxidizer particles decrease to nano-scale dimensions the combustion velocity drastically increases to several km/s [32,33]. These nanocomposite thermites (also termed nanostructured metastable intermolecular composites, nanothermites or superthermites) often have free-standing multilayer foil geometry [34–37]. In contrast to free-standing multilayer thermites in reactive bilayer films coupled with a substrate which is the major heat sink during the reaction, the combustion wave occurs when the temperature of a sample T_S exceeds the initiation temperature T_{in} (T_S > T_{in}) and heating rate is higher than 20 K/s [29,30]. At a heating rate less than 20 K/s the reaction occurs by low reactive diffusion across the entire interface.

In this work, we describe the low-temperature thermite synthesis of Co–ZrO₂ nanocomposite films containing Co nanoclusters separated by ZrO₂ grains. The final goal was to elucidate structural and magnetic transformations formed in the Zr/Co₃O₄ bilayers after thermal annealing from 50 to 500 °C.

2. Experimental section

2.1. Synthesis

The initial films were obtained by the thermal deposition of Co layers at a temperature of 250 °C onto NaCl(001) and chemically pure glass substrates having a thickness of 0.18 mm in vacuum at a residual pressure of 10⁻⁶ Torr. Previously, the substrates were degassed at 350 °C for 1 h. The Co₃O₄ films were obtained by oxidation in an air environment of the initial Co films at a temperature of ~350 °C. The top Zr layer was deposited at room temperature to prevent a reaction between Zr and Co₃O₄ during the deposition. In these experiments, Zr/Co₃O₄ bilayers with an approximate stoichiometry Co:Zr = 3:2 were used. The films on NaCl(001) with the thicknesses of Co equal to 20 nm and Zr - 30 nm were prepared for transmission electron microscopy (TEM) and *in situ* heating electron diffraction studies. In the rest of the experiments films on the glass substrates with a total thickness of 300–400 nm were used. Magnetic, electrical and TEM studies were conducted on Zr/Co₃O₄ bilayers annealed at temperatures between 50 and 500 °C at 50 °C intervals. The samples were held at each temperature for 1 h. All the synthesized Co–ZrO₂ samples were obtained at a pressure of 10⁻⁶ Torr. *In situ* TEM experiments were

performed by heating the Zr/Co₃O₄ films in the column of the JEM-2100 transmission electron microscope (JEOL) using a sample holder with an option of the controlled heating from room temperature to 1000 °C (Gatan Model 652 Double Tilt Heating Holder). The Zr/Co₃O₄ films were separated from the substrate and placed on a molybdenum TEM grid. Electron diffraction data of the *in situ* heating experiments was captured during the experiments.

2.2. Characterization

The phases formed during the synthesis process were identified using a DRON-4007 X-ray diffractometer (Cu K_α radiation). The thicknesses of the Zr and Co₃O₄ layers were determined via X-ray fluorescent analysis. The saturation magnetization M_S, in-plane hysteresis loops and magnetic moment of the samples were measured by a vibrating magnetometer and were also controlled by torque measurements which had been described in detail in our previous work [38]. For finding the Co reduction degree $\eta(T_S)$ versus annealing temperature T_S the magnetic moment M_S^{Co}V₀ of the initial (before oxidation) Co film with V₀ volume was measured. The degree of Co reduction $\eta(T_S)$ was obtained using Formula (2):

$$\eta(T_S) = M_S^{\text{Co}}V(T_S)/M_S^{\text{Co}}V_0 = V(T_S)/V_0 \quad (2)$$

where V(T_S) is Co reduction volume after annealing at temperature T_S.

The resistance R(T_S) measurements were performed by a standard four-probe method in a vacuum of 10⁻⁶ Torr while the as-deposited Zr/Co₃O₄ bilayer annealed at a constant heating rate of 4 °C/min from room temperature up to 500 °C, and then cooling at a rate of ~3 °C/min to room temperature. Magnetoresistance measurements were carried out at room temperature in the magnetic field up to 4.5 kOe by using the standard four-probe technique.

3. Results and discussion

3.1. Synthesis characterization

The initial Co films had saturation magnetization (M_S^{Co} = 1400 emu/cm³) coinciding with the saturation magnetization of the bulk samples. The samples were oxidized in air for the formation of a Co₃O₄ phase which had zero magnetization. To avoid a reaction between Zr and Co₃O₄, the initial Zr/Co₃O₄ bilayers were obtained from the deposition of a Zr layer on the Co₃O₄ film at room temperature. The magnetization M_S of the samples did not change after the deposition of the Zr layer.

Fig. 1 shows the Co reduction degree $\eta(T_S)$ and electrical resistance R(T_S) as a function of annealing temperature T_S for the Zr/Co₃O₄ bilayers. Up to 250 °C the saturation magnetization M_S of the Zr/Co₃O₄ bilayers remained equal to zero ($\eta = 0$), which indicated that there was no mixing and reaction between the Zr and Co₃O₄ layers. With the increasing annealing temperature above 250 °C, the reduction degree η strongly increased up to 300 °C and gradually reached a maximum of $\eta \sim 0.8$ at 500 °C (Fig. 1a). From this it follows that Reaction (1) had the initiation temperature T_{in} ~250 °C and more than 80% of Co was reduced by Zr. Fig. 1b shows the electrical resistance R(T_S) versus annealing temperature T_S for the Zr/Co₃O₄ bilayers which were heated at a rate of 4 °C/min to 500 °C and cooled to room temperature. The temperature increase above T_{in} = 250 °C led to a rapid increase in electrical resistance, which was undoubtedly due to the initiation at T_{in} = 250 °C and development of Reaction (1). The rapid increasing of the electrical resistance was associated with the formation of the insulating ZrO₂ matrix. After annealing from 300 °C to 400 °C the electrical resistance decreased, which was associated with the continuum

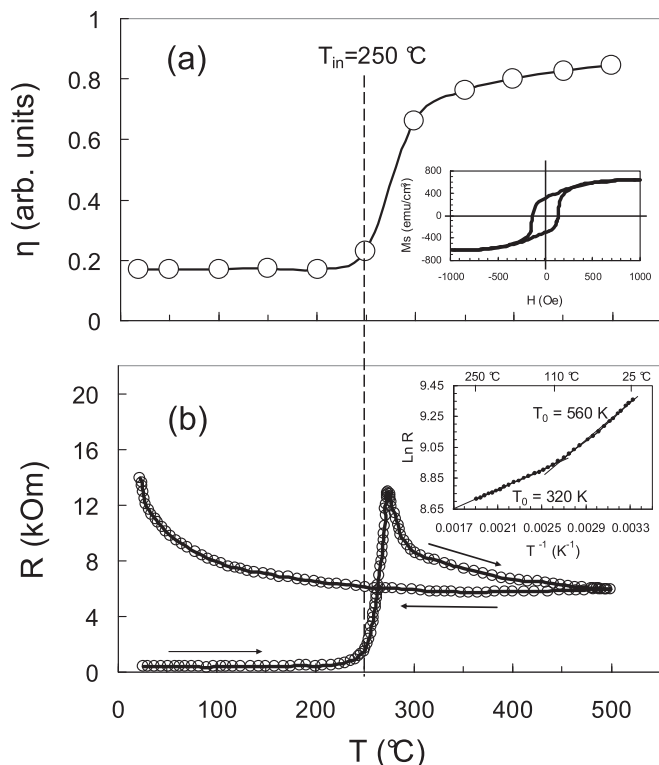


Fig. 1. (a) The degree of Co reduction η as a function of annealing temperature T_s for the Zr/Co₃O₄ bilayers. The inset shows the room temperature hysteresis loop of the synthesized Co–ZrO₂ nanocomposite after annealing at 500 °C. (b) Electrical resistance of the Zr/Co₃O₄ bilayers heated up to 500 °C (the heating rate being 4 °C/min) as a function of temperature measurement T_s . The inset shows the $\ln R$ vs $-T^{-1}$ plot for the as-synthesized Co–ZrO₂ sample after cooling in the 250 °C – 25 °C temperature interval.

reaction (1) and the growth of the Co nanoparticles, which increased the conductivity via an insulating ZrO₂ matrix. Above 400 °C – 500 °C the electrical resistance weakly depended on the annealing temperature and it is assumed that the reaction between Zr and Co₃O₄ had been completed at 500 °C. As the temperature decreased below 500 °C samples began to show metallic-like behavior. The decreasing electrical resistance was most likely associated with the Co grain growth, which eventually formed a percolation cluster in the dielectric ZrO₂ matrix. As the synthesized Co–ZrO₂ samples cooled to the initiation temperature $T_{in} \sim 250$ °C the post-reaction processes continued and promoted the insulation of the Co nanoparticles by ZrO₂ and the start of the increase of the electrical resistance. Upon cooling below 250 °C to room temperature the resistance showed a strong temperature dependence, and it is assumed that the Co–ZrO₂ films were in the insulating conduction regime. For the insulating regime in metal granular films the resistivity R versus temperature T is described by Eq. (3) [1,39–41].

$$R(T) = R_0 \exp[(T_0/T)^n] \quad (3)$$

Where $n = 1$ indicates the thermal activation conduction, $n = 1/2$ tunneling, and $n = 1/4$ Mott variable range hopping. The optimal fitting of the experimental data of the Co–ZrO₂ film has the exponent $n = 1$ for two temperature regions: 20 °C – 110 °C ($T_0 = 560$ K, $E_a^1 = 0.048$ eV) and 110 °C – 250 °C ($T_0 = 560$ K, $E_a^2 = 0.0275$ eV). The $\ln R$ against $-T^{-1}$ plot is given in the inset of Fig. 1b. In both cases the dominant conduction mechanism is thermally activated conduction. One conceivable reason for the low

electrical resistance activation energy is the result of recrystallization and coarsening of Co and ZrO₂ grains, which can occur in the 110 °C – 250 °C temperature interval. After structural relaxation in the as-synthesized Co–ZrO₂ film in the 110 °C – 20 °C temperature interval the as-synthesized Co–ZrO₂ films had an activation energy that is typical for metal granular films at these temperatures [41].

It is significant that only 20% of the Co atoms were not included into the synthesizing metallic Co nanoclusters. Most likely they remained in the unreacted Co₃O₄ and CoO oxides or in the ZrO₂ solid solution, which formed at Co/ZrO₂ interface during Reaction (1) as the transition zone between the Co nanoclusters and ZrO₂ matrix.

3.2. XRD and TEM characterizations

The X-ray diffraction (XRD) pattern of the initial Zr/Co₃O₄ bilayers contained reflections from the fine-dispersed polycrystalline Zr and Co₃O₄ phases (Fig. 2a). After annealing at 300 °C the weak reflections from Co₃O₄ disappeared and a new peak (111) of high-temperature c-ZrO₂ (or t-ZrO₂) phase was formed (Fig. 2b). The c-ZrO₂ and t-ZrO₂ (in the paper c/t-ZrO₂) phases had the proximity of the related peaks and so, these phases are unambiguously distinguished using only conventional X-ray diffraction. The lack of Co peaks implies that in the initial stage Co was fine-crystalline. Upon annealing at 500 °C the X-ray patterns contained (111), (200), (311), (400) c/t-ZrO₂ reflections and (111), (200), (220) reflections of the high-temperature fcc-Co and no peaks from the stable low-temperature hexagonal hcp-Co and monoclinic m-ZrO₂ phases (Fig. 2c). The difference in the binding energies of cobalt in the hexagonal and cubic modifications is insignificant; therefore in thin films and nanostructures the fcc-Co phase is often stabilized.

Fig. 3a shows the sand-like microstructure of the top Zr layer of the initial Zr/Co₃O₄ bilayer, which indicates that the crystalline size of Zr grains is very low (<5 nm). These data are compared with the XRD average crystallite size ~ 5 nm calculated using broadening of the (100)Zr reflection by the Scherrer formula. The interpretation of the electron diffraction reflections (Fig. 3b) confirmed the

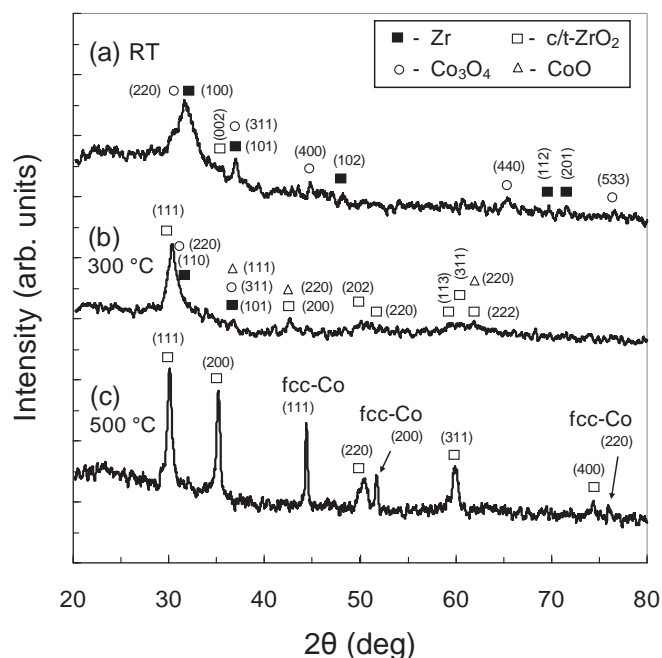


Fig. 2. Typical XRD patterns of the initial Zr/Co₃O₄ bilayers (a) and the synthesized Co–ZrO₂ nanocomposites after annealing at 300 °C (b), 500 °C (c).

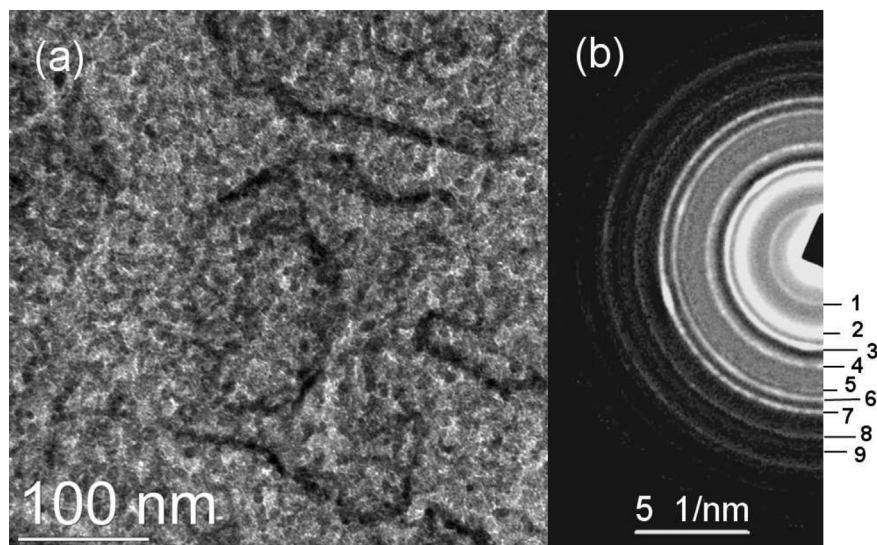


Fig. 3. TEM image (a) and electron diffraction pattern (b) of the initial Zr/Co₃O₄ sample.

coexistence of α -Zr and Co₃O₄ (see Table 1) in the initial Zr/Co₃O₄ films. Fig. 4 shows a typical TEM image, selected area electron diffraction pattern (SAED) of the Co–ZrO₂ film after annealing at 500 °C. The average atomic number for the ZrO₂ phase is lower than the atomic number of Co thus the ZrO₂ region appears brighter on the TEM image (Fig. 4a) than the Co region. The dark regions correspond to the Co grains and the light region - to the ZrO₂ matrix. The TEM observations (Fig. 4a) indicate that the Co nanoparticles have a shape resembling a sphere and were uniformly distributed in the reaction product. The TEM image, SAED pattern (Fig. 4) and electron diffraction analysis suggests that the final reaction product are hcp-Co + fcc-Co nanoparticles embedded in polycrystalline m-ZrO₂ + c/t-ZrO₂ matrix (Table 2). From Fig. 4a it follows that the average diameter of the Co nanoparticles is 20 nm, which is comparable with the superparamagnetic critical size, which is on the order of 20 nm for the Co nanoparticles [42]. However Co nanoparticles with a size below 10 nm are disregarded in the histogram, because it is impossible to obtain a high resolution TEM image due to the film thickness (50 nm). This suggests that all the synthesized Co–ZrO samples contain some amount of superparamagnetic Co nanoparticles. An important point is that the histogram is presented for a film thickness of 50 nm and it may be different from the histograms with film thicknesses of 300–400 nm, which were used for the magnetic, electrical resistance and magnetoresistance investigations.

3.3. Magnetic and magnetoresistance properties

In the Co–ZrO₂ nanocomposites after annealing at 500 °C more

Table 1
Indexing the diffraction reflections in Fig. 3b.

| Rings | Co ₃ O ₄ | α -Zr |
|-------|--------------------------------|--------------|
| 1 | (111) | |
| 2 | (220) | (100) |
| 3 | (311) | (101) |
| 4 | (400) | (102) |
| 5 | (422) | (110) |
| 6 | (511) | (103) |
| 7 | (440) | (112) |
| 8 | (622) | (202) |
| 9 | (711) | |

than 80% of Co was reduced and thus it had high saturation magnetization (~ 500 emu/cm³). The hysteresis loop yielded a typical hysteresis loop of the original Co film (before oxidation) with the coercivity ~ 130 Oe (Fig. 1a). The hysteresis loop achieved saturation at $H_S = 500$ Oe $>$ $H_C = 130$ Oe. This clearly demonstrates the contribution of the superparamagnetic Co particles and disordered spins at Co/ZrO₂ interfaces into the hysteresis loop besides the contribution of the ferromagnetic particles. This result is in good agreement with the results from the TEM analysis, where superparamagnetic Co nanoparticles were assumed. The relatively large value of the ratio M_r to M_S equal 0.62 (where M_r is magnetic remanence) indicates that the Co nanoparticles consisted of randomly oriented equiaxial grains with cubic magnetocrystalline anisotropy [43].

The Co–ZrO₂ samples exhibited negative tunneling magnetoresistance (TMR) which is a characteristic property of granular ferromagnetic metal/insulator nanocomposites. It is known that spin-dependent scattering of the conduction electrons between ferromagnetic nanogranules caused giant TMR in these nanocomposites below the percolation threshold (dielectric regime) [44–47,26]. For the Co–ZrO₂ granular films which had been prepared by pulsed laser deposition the maximum TMR was about 7% at 50 kOe and saturation magnetization was < 50 emu/cm³ [26]. In contrast to these samples the Co–ZrO₂ nanocomposite films which had been prepared by the thermite method revealed very weak negative magnetoresistance about 0.07% at 4.5 kOe and the high saturation magnetization $M_S \sim 500$ emu/cm³ at room temperature.

3.4. In situ electron diffraction study of the Co–ZrO₂ thermite reaction in the Zr/Co₃O₄ thin films

In the initial state the Zr/Co₃O₄ films contained α -Zr and Co₃O₄ layers, as evidenced by the TEM image and corresponding SAED pattern (Fig. 3 and Table 1). The structural changes in the initial Zr/Co₃O₄ bilayer and the formation of the Zr–Co₃O₄ nanocomposite, monitored in the *in situ* heating electron diffraction experiment, is shown in Fig. 5a,b (see the indexing of the diffraction reflections in Table 3a,b). The analysis of the SAED patterns reveals that CoO and amorphous ZrO₂ started forming at 260 °C. This examination proves that the initiation temperature T_{in} of the reaction between Zr and Co₃O₄ was below 260 °C. It is in good agreement with the annealing investigation which indicates that the initiation

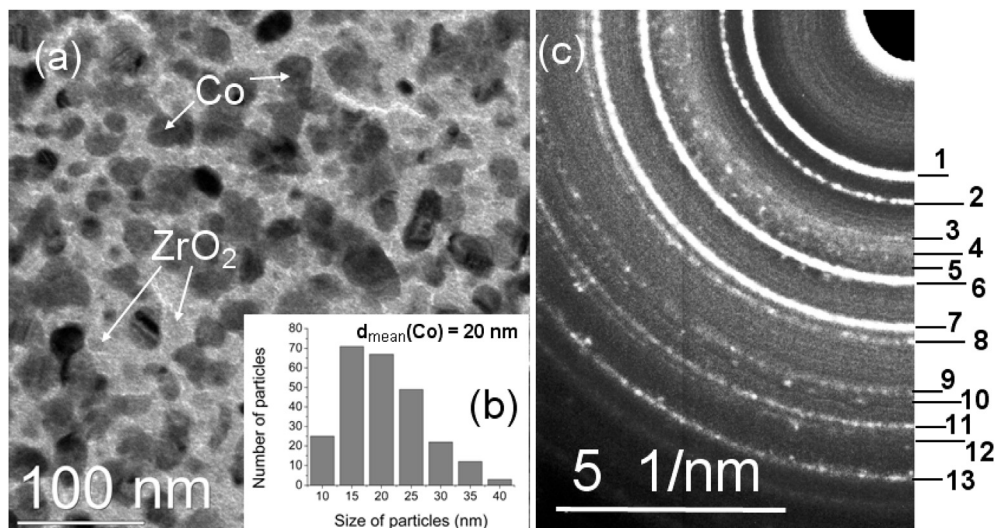


Fig. 4. TEM image (a) of Co nanoparticles embedded in the ZrO₂ matrix, (b) histogram of Co nanoparticles size distribution obtained from the TEM image (a), and (c) electron diffraction pattern of the Co–ZrO₂ film after annealing at 500 °C.

Table 2
Indexing the diffraction reflections in Fig. 4b.

| Rings | m-ZrO ₂ | c/t-ZrO ₂ | Hcp-Co | Fcc-Co |
|-------|--------------------|----------------------|--------|--------|
| 1 | (-111) (111) | (111) | | |
| 2 | (200) | (002) (200) | | |
| 3 | | | (100) | |
| 4 | | | (002) | (111) |
| 5 | | | (101) | |
| 6 | (122) (221) | (202) (220) | | (200) |
| 7 | (213) (312) | (311) | | |
| 8 | (023) (321) | | (102) | |
| 9 | (041) | (004) (400) | (110) | (101) |
| 10 | (-232) | (114) | | |
| 11 | (-313) | (420) | (103) | |
| 12 | (242) | (421) | | (311) |
| 13 | | (511) | (004) | (222) |

Table 3
Indexing the diffraction reflections (a) in Fig. 5a at 300 °C and (b) in Fig. 5b at 500 °C.

| (a) | | | | |
|-------|--------------------|----------------------|--------|--------|
| Rings | m-ZrO ₂ | c/t-ZrO ₂ | Hcp-Co | Fcc-Co |
| 1 | (111) | | | |
| 2 | (-111) | (111) | | |
| 3 | (200) | (002) (200) | | (111) |
| 4 | | | | (200) |
| 5 | (213) (312) | (311) | | (220) |
| 6 | (023) (321) | | | |
| 7 | (041) (004) | (004) (400) | | (311) |
| 8 | (242) | (421) | | (400) |
| 9 | | (511) | | (331) |
| 10 | | | | (422) |
| (b) | | | | |
| Rings | m-ZrO ₂ | c/t-ZrO ₂ | Hcp-Co | Fcc-Co |
| 1 | (111) | | | |
| 2 | (-111) | (111) | | |
| 3 | (200) | (002) | | (111) |
| 4 | | | (002) | |
| 5 | | | (101) | |
| 6 | (122) (221) | (202) (220) | | (200) |
| 7 | (213) (312) | (311) | | |
| 8 | | | (110) | (220) |
| 9 | (314) | (420) | (103) | |
| 10 | | (422) | (004) | (222) |

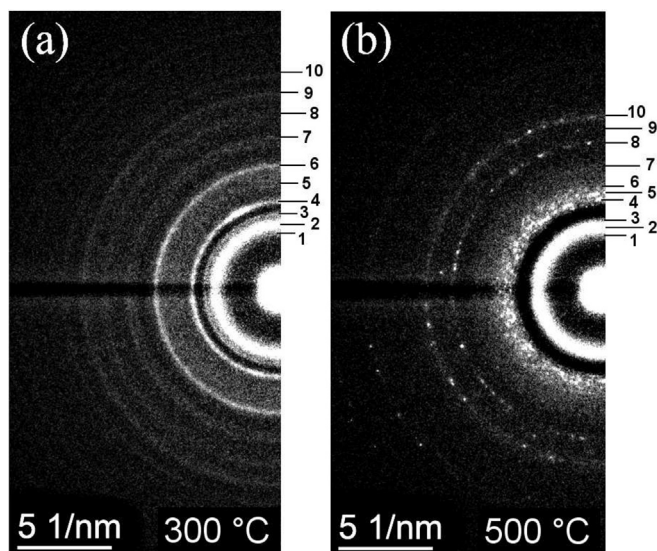


Fig. 5. Electron diffraction patterns of the Co–ZrO₂ film obtained at 300 °C (a) and 500 °C (b).

temperature $T_{\text{in}} \sim 250 \text{ °C}$ (Fig. 1). With increasing the temperature

up to 300 °C the Co₃O₄ reflections gradually disappeared from the SAED patterns and the polycrystalline reflections of the hcp-Co began to form above 400 °C. The existence of fcc-Co cannot be excluded due to the overlap of its main diffraction reflections with those of hcp-Co. The reflections of the CoO phase completely disappeared at 500 °C and the SAED pattern revealed hcp-Co, fcc-Co and m-ZrO₂, c/t-ZrO₂ phases (Fig. 5b and Table 3b). The sample cooling to room temperature did not change the phase composition and final reaction product contained hcp-Co, fcc-Co and m-ZrO₂, c/t-ZrO₂. This result is in good agreement with the final reaction product for the TEM samples obtained by annealing at 500 °C (Fig. 4 and Table 2). While the evidence of the Co formation was apparent from *in situ* heating electron diffraction experiment and XRD studies at temperatures higher 400 °C, the magnetic measurements clearly show the Co formation starting at the initiation temperature

$T_{in} \sim 250$ °C. The most plausible hypothesis is that above $T_{in} \sim 250$ °C Co grains started growing with the fine-dispersed structure which was not detected by the SAED and XRD experiments.

Notice that in the samples with the thickness of 50 nm (prepared for the TEM and *in situ* heating electron diffraction studies) stable hcp-Co and m-ZrO₂ phases were formed, but in the samples for magnetic and electrical investigations with the thickness of above 300–400 nm the high temperature fcc-Co and c/t- ZrO₂ phases, metastable at room temperature, were formed. A possible reason for the formation of the hcp-Co and m-ZrO₂ phases is stress relaxation in the thin layers that created favourable conditions for the equilibrium formation hcp-Co and m-ZrO₂ nanoclusters. However, in the thick layers the reaction stresses insignificantly relaxed and arrested the high temperature fcc-Co and c/t- ZrO₂ phases which were the product of reaction (1).

From the *in situ* heating electron diffraction experiments it follows that the reaction mechanism proposed involves CoO formation in the 260 °C - 500 °C (Table 3a) temperature range and therefore the existence of intermediate reactions (4), (5):



The intermediate reactions often accompany the main reaction in different thermite systems including the classical thermite reaction between Fe₂O₃ and Al (Goldschmidt reaction) [48–50]. Recently we showed that the intermediate reactions were also present in the thermite reaction between Fe₂O₃ and Zr [30].

As is known, thin film solid state reactions are characterized by the initiation temperature T_{in} and the first phase arising in the film interface when the temperature of a sample T_S exceeds T_{in} ($T_S > T_{in}$) [51–54]. We noticed that the initiation temperature $T_{in} \sim 250$ °C of Reaction (1) coincided with the initiation temperature of the thermite reaction in the Zr/Fe₂O₃ thin film [30]. It is common knowledge that the driving force of Al-based thermite reactions is the big negative heat of Al₂O₃ formation ($\Delta H_f = -1676$ kJ/mol). It is reasonable to suggest that the driving force of Zr-based reactions is the negative heat of ZrO₂ formation ($\Delta H_f = -1080$ kJ/mol). Therefore, different reaction systems which are controlled by ZrO₂ synthesis must have the same initiation temperature $T_{in} \sim 250$ °C. Indeed, the value $T_{in} \sim 250$ °C is in good agreement with the initiation temperature $T_{in} \sim 250$ °C in Zr/CuO nanothermite mixture [55] and with the study where the ZrO₂ nanoparticles were obtained by glycothermal processing [56]. In these papers [55,56] using the curves of differential thermal analysis (DTA), thermogravimetric analysis (TGA), and differential scanning calorimetry (DSC), the initiation temperature $T_{in} \sim 250$ °C was defined as the temperature at which heat release started.

Based on all the above, it can be concluded that the final product of the thermite reaction in the Zr/Co₃O₄ thin films contained Co nanoparticles above and below the superparamagnetic critical size embedded into a dielectric ZrO₂ matrix. The particle size of Co during the thermite reaction in the Zr/Co₃O₄ thin films strongly depended on the experimental conditions such as the heating rate and annealing time. The synthesized Co–ZrO₂ samples had good temperature stability and these samples did not change the magnetic properties over 8 months at room temperature. Therefore, the choice of the heating rate, the annealing temperature and annealing time determined the structural and magnetic properties of the nanocomposite Co–ZrO₂ thin films.

4. Summary

In summary, the Co–ZrO₂ ferromagnetic nanocomposites were

synthesized using a new thermite reaction between Zr and Co₃O₄ thin films. The initial Zr/Co₃O₄ films were obtained by the Co films oxidation and sequential deposition of a Zr layers. Thermite reaction (1), using Zr and Co₃O₄ films as a fuel and an oxidizer in Zr/Co₃O₄ bilayers, started above the initiation temperature $T_{in} = 250$ °C. At the annealing temperatures above 300 °C the partial transformation (2) from Co₃O₄ to CoO occurred and thermite reaction (4) between CoO and Zr started. After 500 °C thermite reactions in the Zr/Co₃O₄ films finished and showed a high level of Co reduction (>80%). The final products with a film thickness of 50 nm contained Co nanograins with an average grain size of 20 nm uniformly distributed in the dielectric ZrO₂ matrix. The thermite method provides a promising route to the formation of Co–ZrO₂ ferromagnetic nanocomposite thin films having low initiation temperature, soft magnetic properties, high magnetization and good chemical stability.

Acknowledgments

This study was supported by the Russian Foundation for Basic Research (grants ## 16-03-00069, partially 15-02-00948 and 14-03-00515), by the Council for Grants of the President of the Russian Federation (SP-317.2015.1), and by the program of the Foundation for Assistance to Small Innovative Enterprises in Science and Technology (“UMNIK” program). The TEM studies were carried out using the facilities of Electron Microscopy Laboratory of Siberian Federal University and the Performance Service at Krasnoyarsk Scientific Center.

References

- [1] X. Batlle, A. Labarta, *J. Phys. D: Appl. Phys.* 35 (2002) R15–R42.
- [2] A. Pucci, G. Clavel, M.-G. Willinger, D. Zitoun, N. Pinna, *J. Phys. Chem. C* 113 (2009) 12048–12058.
- [3] T. Wen, K.M. Krishnan, *J. Phys. D: Appl. Phys.* 44 (2011), 393001-1-393001-24.
- [4] J.T. Jiang, X.J. Wei, C.Y. Xu, Z.X. Zhou, L. Zhen, *J. Magn. Magn. Mater.* 334 (2013) 111–118.
- [5] M. Baikousi, O. Kostoula, I. Panagiotopoulos, T. Bakas, A.P. Douvalis, I. Koutselas, A.B. Bourlinos, M.A. Karakassides, *Thin Solid Film.* 520 (2011) 159–165.
- [6] J. Okabayashi, S. Kono, Y. Yamada, K. Nomura, *AIP Adv.* 1 (2011), 042138-1-042138-6.
- [7] A. Butera, J.N. Zhou, J.A. Barnard, *J. Appl. Phys.* 87 (2000) 5627–5629.
- [8] C. Chen, O. Kitakami, S. Okamoto, Y. Shimada, *J. Appl. Phys.* 86 (2161) (1999) 2161–2165.
- [9] J. Gómez, A. Butera, J.A. Barnard, *Phys. Rev. B* 70 (2004), 054428-1-054428-9.
- [10] O. Santini, D.H. Mosca, W.H. Schreiner, R. Marangoni, J.L. Guimaraes, F. Wypych, A.J.A. de Oliveira, *J. Phys. D: Appl. Phys.* 36 (2003) 428–433.
- [11] X. Ni, J. Li, *J. Alloy. Compd.* 558 (2013) 62–67.
- [12] X. Ni, J. Ma, J. Li, D. Jiao, J. Huang, X. Zhang, *J. Alloy. Compd.* 468 (2009) 386–391.
- [13] Y. Chen, A. Li, Y. Li, J. Li, G. Dai, O. Fangping, X. Xiong, *RSC Adv.* 5 (2015) 35929–35933.
- [14] D. Fang, K. Huang, Z. Luo, Y. Wang, S. Liu, Q. Zhang, *J. Mater. Chem.* 21 (2011) 4989–4994.
- [15] P. Manivasakan, V. Rajendran, P. Ranjan Rauta, B.B. Sahu, B.K. Panda, *J. Amer. Ceram. Soc.* 94 (2011) 1410–1420.
- [16] A. Buyukaksoy, T. Fürstenthaupt, V.I. Birss, *Nanoscale* 7 (2015), 8428–8437.
- [17] N. Zink, F. Emmerling, T. Häger, M. Panthöfer, M.N. Tahir, U. Kolb, W. Tremel, *Dalton Trans.* 42 (2013) 432–440.
- [18] S. Ostanin, A. Ernst, L.M. Sandratskii, P. Bruno, M. Däne, I.D. Hughes, J.B. Staunton, W. Hergert, I. Mertig, J. Kudrnovský, *Phys. Rev. Lett.* 98 (2007), 016101-1-016104-4.
- [19] N.H. Hong, M.B. Kanoun, S. Goumri-Said, J.-H. Song, E. Chikoidze, Y. Dumont, A. Ruyter, M. Kurisu, *J. Phys. Condens. Mater* 25 (2013), 436003-1-436003-7.
- [20] A. Pucci, G. Clavel, M.-G. Willinger, D. Zitoun, N. Pinna, *J. Phys. Chem.* 113 (2009) 12048–12058.
- [21] V.V. Kriventsov, D.I. Kochubey, Yu.V. Maximov, I.P. Suzdalev, M.V. Tsodikov, J.A. Navio, M.C. Hidalgo, G. Colón, *Nucl. Instrum. Meth. Phys. Res. A* 470 (2001) 341–346.
- [22] A. Debernardi, D. Sangalli, A. Lamperti, E. Cianci, P. Lupo, F. Casoli, F. Albertini, L. Nasi, R. Ciprian, P. Torelli, *J. Appl. Phys.* 115 (2014), 17D718-1-17D718-3.
- [23] F. Legorreta García, V.G. de Resende, E. De Grave, A. Peigney, A. Barnabé, C. Laurent, *Mater. Res. Bull.* 44 (2009), 1311–1301.
- [24] M. García del Muro, Z. Konstantinović, M. Varela, X. Batlle, A. Labarta, *J. Magn.*

- Magn. Mater 316 (2007) 103–105.
- [25] C. Clavero, B. Sepúlveda, G. Armelles, Z. Konstantinović, M. García del Muro, A. Labarta, X. Batlle, J. Appl. Phys. 100 (2006), 074320-1-074320-5.
- [26] Z. Konstantinović, M. García del Muro, M. Kovylyna, X. Batlle, A. Labarta, Phys. Rev. B 79 (2009), 495304-1-094201-5.
- [27] M. García del Muro, Z. Konstantinović, X. Batlle, A. Labarta, J. Phys. D: Appl. Phys. 46 (2013), 045418-1-045418-9.
- [28] J. Shin, A. Goyal, C. Cantoni, J.W. Sinclair, J.R. Thompson, Nanotechnology 23 (2012), 1555602-1-1555602-6.
- [29] V.G. Myagkov, I.A. Tambasov, O.A. Bayukov, V.S. Zhigalov, L.E. Bykova, Yu.L. Mikhlin, M.N. Volochoev, G.N. Bondarenko, J. Alloy. Compd. 612 (2014) 189–194.
- [30] V.G. Myagkov, L.E. Bykova, O.A. Bayukov, V.S. Zhigalov, I.A. Tambasov, S.M. Zharkov, A.A. Matsynin, G.N. Bondarenko, J. Alloy. Compd. 636 (2015) 223–228.
- [31] L.L. Wang, Z.A. Munir, Y.M. Maximov, J. Mater. Sci. 28 (1993) 3693–3708.
- [32] E.L. Drezin, Prog. Energy Combust. Sci. 35 (2009) 141–167.
- [33] Y. Yang, D.-r. Yan, Y.-c. Dong, X.-g. Chen, L. Wang, Z.-h. Chu, J.-x. Zhang, J.-n. He, J. Alloy. Compd. 579 (2013) 1.
- [34] K.J. Blobaum, M.E. Reiss, J.M. Plitzko, T.P. Weihs, J. Appl. Phys. 94 (2003) 2915–2922.
- [35] P. Zhu, J. Jiao, R. Shen, Y. Ye, Shuai Fu, D. Li, J. Appl. Phys. 115 (2014), 194502-1-194502-5.
- [36] E.J. Mily, A. Oni, J.M. LeBeau, Y. Liu, H.J. Brown-Shaklee, J.F. Ihlefeld, J.-P. Maria, Thin Solid Film. 562 (2014) 405–410.
- [37] M. Petrantonì, C. Rossi, L. Salvagnac, V. Conédéra, A. Estève, C. Tenaillieu, P. Alphonse, Y.J. Chabal, J. Appl. Phys. 108 (2010), 084323-1-084323-5.
- [38] V.G. Myagkov, V.S. Zhigalov, L.E. Bykova, G.N. Bondarenko, J. Magn. Mater. 305 (2006) 534–545.
- [39] B.I. Shklovskii, A.L. Efros, Electronic Properties of Doped Semiconductors, Springer Verlag, Berlin, 1984.
- [40] N.F. Mott, Conduction in Non-crystalline Materials, Oxford University Press, Oxford, 1987.
- [41] C.A. Neugebauer, M.B. Webb, J. Appl. Phys. 33 (1962) 74–82.
- [42] A.I. Gusev, A.A. Rempel, Nanocrystalline Materials, Cambridge International Science Publishing, 2004, p. p.351.
- [43] J.H. Fendle (Ed.), Nanoparticles and Nanostructured Films: Preparation, Characterization and Applications, WILEY_VCH Verlag GmbH & Co KGaA, NY, 1998.
- [44] S. Mitani, H. Fujimori, K. Takanashi, K. Yakushiji, J.-G. Ha, S. Takahashi, S. Maekawa, S. Ohnuma, N. Kobayashi, T. Masumoto, M. Ohnuma, K. Hono, J. Magn. Mater. 198–199 (1999) 179–184.
- [45] A. Milner, A. Gerber, B. Groisman, M. Karpovsky, A. Gladkikh, Phys. Rev. Lett. 76 (1996) 475–478.
- [46] S. Sankar, A.E. Berkowitz, D.J. Smith, Phys. Rev. B 62 (2000) 14273–14278.
- [47] A. García-García, J.A. Pardo, P. Štrichovanec, C. Magén, A. Vovk, J.M. De Teresa, G.N. Kakazei, Y.G. Pogorelov, L. Morellón, P.A. Algarabel, M.R. Ibarra, Appl. Phys. Lett. 98 (2011), 122502-1-122502-3.
- [48] Y. Yang, Y.-c. Dong, X.-g. Chen, L. Wang, Z.-h. Chu, J.-x. Zhang, J.-n. He, J. Alloy. Compd. 579 (2013) 1–6.
- [49] A. Babakhani, E. Zahabi, H.Y. Mehrabani, J. Alloy. Compd. 514 (2012) 20–24.
- [50] J.L. Chen, H.H. Hng, Y.W. Lee, S.W. Du, N.N. Thadhani, Combust. Flame 157 (2010) 2241–2249.
- [51] J.M. Poate, K.N. Tu, J.W. Mayer (Eds.), Thin Films-interdiffusion and Reaction, Wiley-Interscience, New York, 1978.
- [52] E.G. Colgan, Mater. Sci. Rep. 5 (1990) 1–44.
- [53] R. Pretorius, C.C. Theron, A. Vantomme, J.W. Mayer, Crit. Rev. Solid. State Mater. Sci. 24 (1999) 1–62.
- [54] T. Laurila, J. Molarius, Crit. Rev. Solid. State Mater. Sci. 28 (2003) 185–230.
- [55] E.J. Mily, A. Oni, J.M. LeBeau, Y. Liu, H.J. Brown-Shaklee, J.F. Ihlefeld, J.-P. Maria, Thin Solid Film. 562 (2014) 405–410.
- [56] J.-S. Kim, D.-H. Lee, S. Kang, D.-S. Bae, H.-Y. Park, M.-K. Na, Trans. Nonferrous Met. Soc. China 19 (2009) s88–s91.