

# Using Redox-Active $\pi$ Bridging Ligand as a Control Switch of Intramolecular Magnetic Interactions

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## Supporting Information

**ABSTRACT:** Intramolecular magnetic interactions in the dinuclear complexes  $[(\text{tpy})\text{Ni}(\text{tphz})\text{Ni}(\text{tpy})]^{n+}$  ( $n = 4, 3,$  and  $2$ ; tpy, terpyridine; tphz, tetrapyridophenazine) were tailored by changing the oxidation state of the pyrazine-based bridging ligand. While its neutral form mediates a weak antiferromagnetic (AF) coupling between the two  $S = 1$  Ni(II), its reduced form,  $\text{tphz}^{\bullet-}$ , promotes a remarkably large ferromagnetic exchange of  $+214(5)$  K with Ni(II) spins. Reducing twice the bridging ligand affords weak Ni–Ni interactions, in marked contrast to the Co(II) analogue. Those experimental results, supported by a careful examination of the involved orbitals, provide a clear understanding of the factors which govern strength and sign of the magnetic exchange through an aromatic bridging ligand, a prerequisite for the rational design of strongly coupled molecular systems and high  $T_C$  molecule-based magnets.

Efficient control of the interactions between magnetic centers is a fascinating challenge, at the basis of any rational design of molecule-based magnets.<sup>1</sup> This fundamental problem is intrinsically linked to the question of the relative localization (or delocalization) of the unpaired electrons. Those are often located on  $d$  (or  $f$ ) orbitals of metal ions, and they interact weakly by superexchange through diamagnetic bridging ligands or atoms. Nevertheless, alternative strategies have recently received an increasing interest to promote strong magnetic coupling: (i) the mixed-valence approach, for which an electron delocalization promotes double-exchange, which exceeds superexchange;<sup>2</sup> and (ii) the radical bridge approach using an organic radical ligand acting as a magnetic relay between paramagnetic metal centers.<sup>3</sup> For the latter, the organic bridge characteristics are crucial to optimize the strength of the magnetic interactions: it should be strongly coordinating for an efficient mixing of the involved metal/organic orbitals,<sup>3d,4</sup> and its unpaired electron as delocalized as possible between the paramagnetic metal ions.

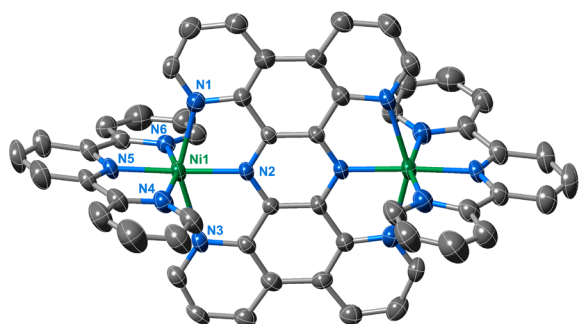
Based on these prerequisites, our group reported recently<sup>4</sup> a prototype dinuclear complex based on Co(II) and the redox-active tphz bridging ligand.<sup>5</sup> In this system, the intramolecular magnetic interactions are controlled by successive reduction

processes. While the neutral form of tphz mediates a weak AF exchange between  $ls$ -Co(II) ( $ls$ : low-spin), the once reduced radical form stabilizes very strong AF couplings between the  $S = 1/2$  organic spin and the two  $hs$ -Co(II) ( $hs$ : high-spin). The resulting high-spin complex possesses a well isolated  $S_T = 5/2$  ground state that is remarkably the only thermally populated state at ambient temperature. When further reducing this dinuclear complex, despite the formal diamagnetic state of the  $\text{tphz}^{2-}$  ligand, a strong AF coupling between two  $hs$ -Co(II) is observed, leading to an overall nonmagnetic ground state. In this case, the strong magnetic interaction is induced by a large spin delocalization arising from an efficient orbital mixing of the  $\text{tphz}^{2-}$   $\pi$  system and one of the Co(II) singly occupied  $d$  orbitals. Therefore, in this type of radical/metal ion complexes, the nature and magnitude of the intramolecular coupling is indeed predictable based on both electronic configurations of metal ions and the bridging ligand.

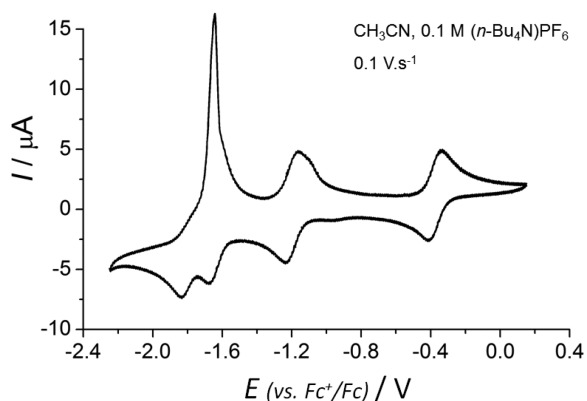
To further experimentally investigate the role of the metal/organic orbital complementarity in the control of the magnetic interactions, a novel family of dinuclear Ni(II) complexes was synthesized with the redox-active tphz ligand. The reaction of  $\text{Ni}(\text{tpy})\text{Cl}_2$ ,<sup>6</sup>  $\text{TiPF}_6$  and tphz in  $\text{CH}_3\text{CN}$  leads to a new dinuclear nickel complex that was isolated as orange needle-shaped crystals of  $[\text{Ni}_2(\text{tphz})(\text{tpy})_2](\text{PF}_6)_4 \cdot 3\text{CH}_3\text{CN}$  (**1**) after slow diffusion of  $\text{Et}_2\text{O}$  vapors (Supporting Information, SI). The single crystal X-ray diffraction data at 120 K (Figure 1) reveal that **1** is isostructural to the Co(II) analogue,<sup>5</sup> possessing two equivalent  $\text{Ni}^{\text{II}}(\text{tpy})^{2+}$  fragments bridged by a neutral tphz ligand. The cyclic voltammetry of **1** in  $\text{CH}_3\text{CN}$  reveals four one-electron redox events at  $-0.40$ ,  $-1.25$ ,  $-1.66$ , and  $-1.82$  V versus vs  $\text{Fc}^+/\text{Fc}$  (Figure 2)<sup>7</sup> with a rest potential at  $-0.15$  V, indicating only reduction processes. By comparing with the reported Ru(II)<sup>8</sup> and Co(II)<sup>4</sup> analogues, the two processes at  $-0.40$  and  $-1.25$  V can be assigned to the reduction of the bridging ligand and the other two most cathodic ones to the  $\text{tpy}$  reduction. The large separation between the two first redox potentials evidence the high stability of the tphz radical in the complex, with a comproportionation constant of ca.  $2.6 \times 10^{14}$ . Based on these redox properties, **1** was chemically reduced with  $\text{KC}_8$  to

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**Figure 1.** Crystal structure of  $[\text{Ni}_2(\text{tphz})(\text{tpy})_2]^{4+}$  in **1** at 120 K (thermal ellipsoids: 50%). Hydrogen atoms, solvent molecules and anions are omitted for clarity.



**Figure 2.** Cyclic voltammogram for a solution of **1**.

give once and twice-reduced systems as  $[\text{Ni}_2(\text{tphz})(\text{tpy})_2](\text{PF}_6)_3 \cdot 2\text{Et}_2\text{O}$  (**2**) and  $[\text{Ni}_2(\text{tphz})(\text{tpy})_2](\text{PF}_6)_2 \cdot \text{Et}_2\text{O}$  (**3**) (SI) with similar coordination environments in the dinuclear cationic Ni(II) complexes (Figure S1 and Table S1). Nevertheless, significant differences are seen when comparing the bond distances within the bridging ligand (Tables 1 and S2) in agreement with two successive reductions as observed for the Co series.<sup>4</sup> In particular, C–N and N2...N2' distances in the tphz pyrazine ring are significantly elongated, highlighting the successive reductions mainly centered on the bridging ligand. The Ni1–N2 distance decreases from 1.980(4) to 1.921(4) Å upon the two reductions as a result of the increasing electrostatic interactions between the cationic metal and the bridging ligand. Those bond distances are particularly short, but in perfect agreement with the reported cobalt series.<sup>4</sup> The average Ni–N bond distance involving terpyridine remains quasi-unchanged upon these two first

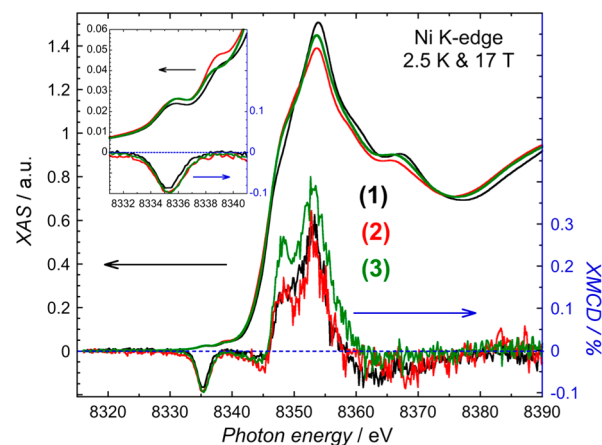
**Table 1.** Selected Bond Distances (Å) in **1–3**

	<b>1</b>	<b>2</b>	<b>3</b>
Ni–N1	2.268(5)	2.267(7)	2.305(4)
Ni–N2	1.980(4)	1.962(6)	1.921(4)
Ni–N3	2.336(5)	2.338(7)	2.326(4)
Ni–N4	2.089(5)	2.097(7)	2.100(4)
Ni–N5	1.975(4)	2.002(6)	1.997(3)
Ni–N6	2.078(5)	2.084(7)	2.114(4)
Ni–N(tpy) <sup>a</sup>	2.047(5)	2.061(7)	2.070(4)
C–C(pz) <sup>a</sup>	1.391(7)	1.384(10)	1.382(6)
C–N(pz) <sup>ab</sup>	1.331(7)	1.359(9)	1.375(6)

<sup>a</sup>pz: pyrazine part of tphz. <sup>b</sup>Average distance.

reductions, indicating the tpy redox innocence and that the Ni ion preserved its spin state and charge in agreement with CASSCF calculations (Figure S4).<sup>9</sup>

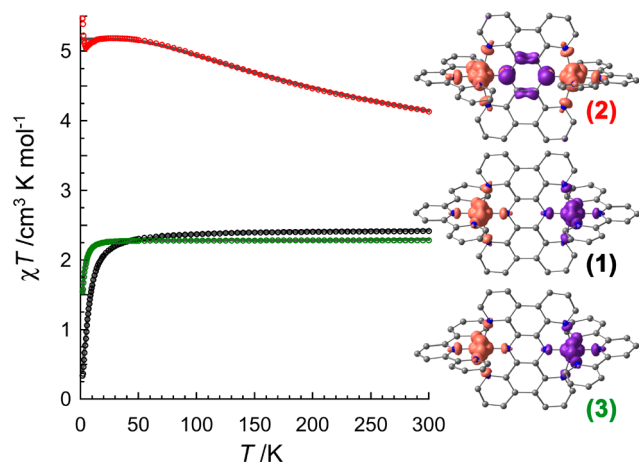
The local electronic and magnetic properties of the nickel metal ions were studied by X-ray absorption spectroscopy (XAS) and X-ray magnetic circular dichroism (XMCD). These techniques are particularly relevant for complexes with redox-active ligands, for which the oxidation state of the metal ion can be questioned.<sup>10,11</sup> The XAS spectrum at the Ni K-edge (Figure 3) is dominated by the  $1s \rightarrow 4p$  transitions which show



**Figure 3.** XAS and XMCD spectra. XAS spectra were normalized to zero before the edge and to unity far above the edge. XMCD spectra are given in XAS spectra %. Inset: magnification of the pre-edge region.

similar maxima for the three complexes between 8353.5 and 8353.8 eV. As expected for comparable coordination sphere of the probed atom, the first EXAFS (Extended X-ray Absorption Fine Structure) oscillations overlap remarkably well. The much weaker pre-edge, dipole-forbidden,  $1s \rightarrow 3d$  transitions are detected at  $8335.9 \pm 0.2$  and  $8338.9 \pm 0.3$  eV for the three complexes in agreement with the expected and calculated Ni(II) ligand field splitting scheme (Figure S4).<sup>11a</sup> The XMCD experiments were also performed at the nickel K-edge to probe the Ni orbital magnetic moment (Figure 3). Considering that the spin–orbit interactions of the ligand atoms are negligibly small and thus cannot induce a significant orbital magnetic moment on the Ni centers, the observed XMCD signal at  $8335.2 \pm 0.1$  eV is only due to the orbital polarization of the  $4p$  and  $3d$  states induced by the intra-atomic spin–orbit coupling of the Ni atoms. The Ni magnetic moment is thus directly proportional to the XMCD signals observed in the pre-edge and edge regions. Their quasi-identical intensities demonstrate the presence of the same magnetic moment and thus spin-state on the Ni metal ions in **1–3**. The similarity of the XAS and XMCD spectra establishes unambiguously the identical electronic and magnetic characteristics of the Ni site in these complexes.

To quantify the strength of the intramolecular magnetic coupling, dc magnetic measurements were performed at 0.1 T (Figure 4). The  $\chi T$  product of **1** remains constant between 300 and 80 K at  $2.4 \text{ cm}^3 \text{ K mol}^{-1}$ , in good agreement with the presence of two  $S = 1$  Ni(II) ( $C_{\text{Ni}} = 1.2 \text{ cm}^3 \text{ K mol}^{-1}$ ;  $g_{\text{Ni}} = 2.19(5)$ ). Below 80 K, the progressive decrease of  $\chi T$ , to a minimum value of  $0.38 \text{ cm}^3 \text{ K mol}^{-1}$  at 1.85 K, is the result of the concerted effect of weak AF interactions between nickel



**Figure 4.**  $\chi T$  vs  $T$  plot at 0.1 T (gray lines: best fits; see text). Insets: View of the computed spin density.

spins and the Ni magnetic anisotropy. The fit of the experimental data with both contributions (solid lines, Figures 4 and S8)<sup>12,13</sup> gives  $g_{\text{Ni}} = 2.21(5)$ ,  $D_{\text{Ni}}/k_{\text{B}} = +13.6(3)$  K and  $J_{\text{Ni-Ni}}/k_{\text{B}} = -1.86(3)$  K comparable to  $-4.4$  K obtained by broken-symmetry DFT (B3LYP/def2-TZVP) calculations (See SI for details).<sup>9</sup>

For the once-reduced analogue, **2**, the  $\chi T$  product at 300 K is  $4.1 \text{ cm}^3 \text{ K mol}^{-1}$ , significantly higher than ca.  $2.8 \text{ cm}^3 \text{ K mol}^{-1}$  expected for one  $S = 1/2$  radical ( $g_{\text{rad}} = 2.0$ ) and two Ni(II) ( $g_{\text{Ni}} = 2.2$ ,  $S = 1$  in agreement with XAS/XMCD). This result suggests a remarkably strong ferromagnetic exchange between the metal ion and the bridging radical  $\text{tphz}^{\bullet}$  spins that is confirmed by the  $\chi T$  increase between 300 and 40 K (Figures 4 and S9). Below 40 K, a  $\chi T$  plateau is reached with a maximum value of  $5.2 \text{ cm}^3 \text{ K mol}^{-1}$ , which corresponds to an  $S_{\text{T}} = 5/2$  ground state ( $g_{S_{\text{T}}=5/2} = 2.18(5)$ ). In this system, the ferromagnetic coupling is straightforward to rationalize considering the orthogonality of the Ni magnetic orbitals ( $d_{x^2-y^2}$  and  $d_{z^2}$ )<sup>14</sup> and the bridging ligand SOMO which is delocalized on the  $\pi$  system. Using an isotropic spin Heisenberg model,<sup>15</sup> its magnitude was estimated<sup>13</sup> to be  $J_{\text{Ni-rad}}/k_{\text{B}} = +214(5)$  K ( $g = 2.17(5)$ ; Figure S9) leading to an energetically well isolated  $S_{\text{T}} = 5/2$  ground state (the first  $S = 3/2$  excited state lies at 214 K above). This large coupling is in good agreement with the calculated one (+249 K; DFT; SI),<sup>9</sup> and stronger than what was observed in other radical bridged Ni(II) complexes,<sup>16</sup> likely as a consequence of the strong metal–ligand coordination revealed by the short Ni–N2 bond.

For the twice-reduced compound **3**, the magnetic susceptibility data are similar to that for **1** with a constant  $\chi T$  product at ca.  $2.3 \text{ cm}^3 \text{ K mol}^{-1}$  between 300 and 25 K (Figures 4 and S10), which is expected for a Curie behavior with two Ni(II) spins ( $g_{\text{Ni}} = 2.14(5)$ ). As suggested by structural and spectroscopic measurements, the magnetic properties conclude unambiguously that the Ni oxidation and spin states do not change after successive reductions. Upon cooling below 25 K, the  $\chi T$  value decreases to  $1.5 \text{ cm}^3 \text{ K mol}^{-1}$  at 1.85 K. Like in **1**, this low temperature behavior is probably the result of the combined contributions from the weak intramolecular AF interaction ( $-2.2$  K from DFT)<sup>9</sup> and the Ni magnetic anisotropy. The  $\chi T$  vs  $T$  and  $M$  vs  $H$  fits<sup>12,13</sup> lead systematically to  $J_{\text{Ni-Ni}}$  values close to zero ( $|J_{\text{Ni-Ni}}/k_{\text{B}}| < 0.1$  K) with

$g_{\text{Ni}} = 2.14(5)$  and  $D_{\text{Ni}}/k_{\text{B}} = -10.4(3)$  K (solid lines in Figures 4 and S10). Despite the presence of a significant magnetic anisotropy, no out-of-phase ac signal was detected above 1.85 K and up to 10 kHz for these compounds.

This new series of complexes with different redox states offers the appealing opportunity for a direct comparison between Ni(II) and Co(II) analogues, for which only the metal electronic configuration differs. Indeed, the once- and twice-reduced forms of the Ni and Co complexes are quasi-isometric, providing a unique platform to directly probe the influence of the singly occupied atomic orbitals of the metal ion on the strength and sign of the magnetic coupling. According to the ab initio ligand field theory, the Ni(II) unpaired electrons are formally located on the  $d_{x^2-y^2}$  and  $d_{z^2}$  orbitals (Figure S4),<sup>14</sup> which are only weakly interacting with  $\text{tphz}$  nitrogen  $\sigma$ -orbitals (Figures S5 and S7). In this situation, the spin density is, as expected, mostly located on the metal ions in **1** and **3** (Figure 4) and the magnetic interactions between the two Ni(II) spins are necessarily weak. Moreover, those metal ion singly occupied orbitals are orthogonal to the ligand  $\pi$  system. Therefore, as already mentioned, when the  $\text{tphz}^{\bullet}$  radical is stabilized in **2**, the unpaired electron is located on the ligand  $\pi$  orbital and a large ferromagnetic coupling is observed between metal and radical spins.

In  $hs$ -Co(II) analogues, in addition to the two unpaired electrons located on the  $d_{x^2-y^2}$  and  $d_{z^2}$  orbitals, a third one is available on  $d_{xy}$ . This different configuration has a considerable impact on the magnetism,<sup>4</sup> as  $d_{xy}$  is the only orbital with an orientation that allows a significant overlap with the ligand  $\pi$  system. Hence, this efficient mixing of the  $d_{xy}$  and  $\pi$  orbitals induces an efficient spin density delocalization in the  $hs$ -Co(II) complexes. When twice-reduced,  $\text{tphz}^{2-}$  is formally diamagnetic, but the large spin delocalization promotes a strong AF coupling between the two Co(II) ( $-74$  K), in striking contrast to **3**.<sup>17</sup> This comparative study demonstrates the key role of the  $d_{xy}$  orbital on the strength of the intramolecular magnetic exchange.

At least one order of magnitude is also gained on the magnetic interaction, when a single electron is located on the ligand  $\pi$  orbital. The metal ion and  $\text{tphz}^{\bullet}$  spins are strongly ferromagnetically coupled in the Ni(II) case with  $J/k_{\text{B}} = +214$  K (vide supra), while a huge antiferromagnetic exchange is observed for the  $hs$ -Co(II) complex (at least  $-500$  K). The  $J(d_{xy}/\pi)$  coupling between the  $d_{xy}$  and  $\pi$  unpaired electrons is thus very large as it overcomes the ferromagnetic  $J(d_{x^2-y^2}/\pi)$  and  $J(d_{z^2}/\pi)$  contributions. Qualitatively, the  $J(d_{xy}/\pi)$  interaction is thus larger than  $-700$  K from the simple  $J_{\text{Co-rad}}/k_{\text{B}} - J_{\text{Ni-rad}}/k_{\text{B}}$  relation.<sup>18</sup>

In conclusion, this work reports on a new series of dinuclear  $[(\text{tphz})\text{Ni}(\text{tphz})\text{Ni}(\text{tphz})]^{n+}$  complexes which provides, in comparison to the  $hs$ -Co(II) analogues, a complete set of experimental and theoretical results allowing a general understanding of the dominant factors governing the strength and nature of the magnetic coupling via an aromatic bridge. In particular, in addition to the obvious interest of using a radical bridging ligand, the careful selection of the metal ion based on its electronic configuration is the key step to promote the strongest possible coupling; this is happening when the unpaired electrons reside on the suitable  $t_{2g}$  orbitals, which overlap the best with the radical  $\pi$  system. This general

approach should be easily transposed to systems with higher nuclearities and dimensionalities, for which strong magnetic exchanges could be fully exploited to design high temperature molecule-based magnets.<sup>10,19</sup>

## ■ ASSOCIATED CONTENT

### 📄 Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/jacs.9b03044.

Experimental details about syntheses, structural, spectroscopic, magnetic characterizations and theoretical calculations (PDF)

X-ray crystallographic data for **1** (CIF)

X-ray crystallographic data for **2** (CIF)

X-ray crystallographic data for **3** (CIF)

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### Notes

The authors declare no competing financial interest.

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$$\hat{H} = -2J_{\text{Ni-Ni}}(\vec{S}_{\text{Ni1}} \cdot \vec{S}_{\text{Ni2}}) + 2D_{\text{Ni}}S_{\text{Ni,z}}^2 + 2g_{\text{Ni}}\mu_{\text{B}}\vec{S}_{\text{Ni}} \cdot \vec{H}$$

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(14) The average N1–Ni–N3 axis is taken as the z direction by convention (Figure 1).

(15) The following Hamiltonian has been used to fit the magnetic data:

$$\hat{H} = -2J_{\text{Ni-rad}}(\vec{S}_{\text{Ni1}} + \vec{S}_{\text{Ni2}}) \cdot \vec{S}_{\text{rad}} + \mu_{\text{B}}(2g_{\text{Ni}}\vec{S}_{\text{Ni}} + g_{\text{rad}}\vec{S}_{\text{rad}}) \cdot \vec{H}$$

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(18) More accurately, this  $J(d_{xy}/\pi)$  antiferromagnetic coupling is better described as being between the  $d_{xy}(\text{Co}) + \pi(\text{N}_{\text{tphz}})$  orbital and the  $\pi$  system of the carbon atoms of the tphz pyrazine part (for symmetry reasons of the involved orbitals).

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