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Solid-state synthesis, magnetic and structural properties of epitaxial $D0_3$ -Fe₃Rh(001) thin films

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Here we first report on the formation of epitaxial D0₃-Fe₃Rh(001) films grown during solid-state reaction in Rh/Fe(001) bilayers on MgO(001) substrates. Samples of the Fe₆₈Rh₃₂ composition above 400 °C showed the formation of a new ordered phase, in addition to the ordered B2–FeRh phase (space group Pm-3m, lattice constant a = 0.2993 nm), which becomes the dominant phase in Fe₇₆Rh₂₄ samples. These results and the results of the asymmetrical XRD φ -scan prove that the new ordered phase is the ordered D0₃-Fe₃Rh(001) phase (space space spa

1. Introduction

Magnetic anisotropy

Among Fe_xRh_{100-x} alloys, the alloys close to equatomic ($48 \le x \le 56$) composition in both bulk samples and thin films have attracted a great deal of interest in recent years due to their unique first-order metamagnetic transition from an antiferromagnetic α "-phase to a ferromagnetic α '-phase when heated above ~ 370 K [1]. An unusual feature of the reversible $\alpha'' \leftrightarrow \alpha'$ transition is that both phases have an ordered B2 structure and a slight 1% isotropic difference in the lattice constant, which gives existence to the giant magnetostriction [2], large magnetoresistance [3,4] and magnetocaloric effects [5]. This expands the possible use of B2-FeRh thin films for various applications such as heat assisted magnetic recording (HAMR) [6-8], micro-electro-mechanical systems (MEMS) [9] and spintronic devices [10-12]. According to the bulk Fe-Rh phase diagram, the FexRh100-x alloys are not magnetic and have an fcc $\gamma(A1)$ structure at room temperature when Fe is less than 48 at.% [13,14]. The chemically ordered bcc (B2) structure exists in a wide composition range (48 \leq x \leq 85) and at compositions of (85 \leq x \leq 100) the alloys are in a chemically disordered bcc (A2) structure [13-16]. Comprehensive high-performance first-principles calculations based on density functional theory predict stable ordered structures of Fe₃Rh (P4/mmm), Fe₂Rh (C11_b) and FeRh₃ (D0₂₄) in the binary Fe-Rh system

[17,18] in addition to B2–FeRh. The magnetic moment and anisotropy of Rh_2Fe_m ($m \le 4$) and ordered $Fe_2Rh(C11_b)$ structures were determined using first-principle calculations based on the Vienna ab initio simulation package in ref. 19. The calculations predict the Rh atoms are ferromagnetically coupled to the Fe, the interface Fe moments are about 3.0 µB for Rh₂Fe_m and C11_b-Fe₂Rh alloys and the largest uniaxial anisotropy is found in Fe₂Rh ($7 \cdot 10^7$ erg/cm³) [19]. However, the currently known experimental results are often incompatible with theoretical predictions of new phases. It follows from experimental data that, apart from the chemically ordered B2-FeRh phase, there are no other ordered phases in the Fe-Rh system. For many binary alloys, the compositional region of the B2 phase existence is centered near the equiatomic composition. However, in the Fe-Rh system this region is strongly shifted to a region with a high Fe concentration, which suggests the existence of an Fe-rich phase with lattice parameters coinciding or very close to B2-FeRh Back in 1963 G. Shirane et al. [13] expected a DO_3 ordering in the Fe-rich part of the Fe-Rh binary. However, efforts to obtain the Fe₃Al type (D0₃ structure) ordering in 25 at.% rhodium and 75 at.% iron alloys by cooling from 550 to 300 °C at a rate of 1 °C/h was not successful [13]. Despite numerous studies over the past 60 years, the binary Fe-Rh phase diagram outside the (48 \leq x \leq 56) iron composition is still being refined [13,14].

group Fm-3m, lattice constant a = 0.5888 nm), forming a cube-on-cube orientation relationship with respect to

the MgO(001) substrate. The $D0_3$ -Fe₃Rh sample is a soft magnetic material with high saturation magnetization.

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In this paper, we report the investigations of solid-state reaction between polycrystalline Rh and epitaxial Fe(001) films in 32Rh/68Fe (001) and 24Rh/76Fe(001) bilayers with Fe₆₈Rh₃₂ and Fe₇₆Rh₂₄ atomic ratios, respectively. In both samples, a mixture of epitaxially intergrown D0₃-Fe₃Rh(001) and B2–FeRh(001) phases is formed above 400 °C. The lattice parameter, saturation magnetization and the first constant of magnetocrystallographic anisotropy K_1 of the D0₃-Fe₃Rh phase were determined. A modified Fe-rich part of the Fe–Rh phase diagram was proposed based on the results obtained in this study.

2. Experimental

Iron-rich 32Rh/68Fe(001) bilayers grown on a MgO(001) substrate were fabricated in two steps. In the first step, the epitaxial Fe(001)/MgO (001) films were grown on single-crystal MgO(001) substrates with a size of $10 \times 10 \times 1$ mm by a thermal evaporation method in a vacuum chamber at a pressure of $1.3 \cdot 10^{-4}$ Pa. To obtain high-quality Fe(001) films, the MgO(001) substrates were previously outgassed at 300 °C for 1 h and the Fe layers were deposited at 250 °C. Epitaxial Fe(001) films had the orientation ratio Fe(001)[100] || MgO(001)[110] with the MgO (001) substrate.

At the second step, a Rh layer was deposited onto the Fe(001)/MgO (001) samples by DC magnetron sputtering. The base pressure of the vacuum chamber was $\sim 1.3 \cdot 10^{-4}$ Pa, the operating pressure was $\sim 1.3 \cdot 10^{-1}$ Pa, in an Ar atmosphere. To prevent a reaction between Rh and Fe, the Rh layer was deposited on Fe(001) at room temperature. Under such deposition conditions the polycrystalline Rh layer was formed on the Fe (001) surface. The thickness of the Fe(001) layer was 150 nm, and the thicknesses of the upper Rh layer was 80 nm for the Rh/Fe(001) bilayers.

To investigate the total elemental composition of the films, analysis was performed by energy-dispersive X-ray spectroscopy (EDS) using a Hitachi TM4000Plus instrument (Tokyo, Japan) equipped with an EDS detector. EDS analysis was performed in a high vacuum at an acceleration voltage of 20 kV with $100 \times$ magnification from a film area of 1.5 mm². EDS analysis performed for the initial Rh/Fe(001) bilayers showed the elemental composition of 68 at.% Fe and 32 at.% Rh (Fe68Rh32).

The initial samples were annealed under a vacuum of $1.3 \cdot 10^{-4}$ Pa from room temperature to 800 °C in intervals of 100 °C, holding for 1 h at each temperature. The phase formations and structural changes occurring at the interface between the Rh and Fe layers when the annealing temperature increased, were identified with a DRON-4-07 diffractometer (CuK_{α} radiation, $\lambda = 0.15406$ nm). The epitaxial film orientations were analyzed by means of asymmetric X-ray diffractometer (Almelo, The Netherlands) with a PIXcel detector with a PIXcel detector. CuK_{α} ($\lambda = 0.15406$ nm) radiation monochromatized by a secondary graphite monochromator was used in the instrument. The XRD investigations were carried out using the samples at room temperature.

The saturation magnetization $M_{\rm S}$ and the coercivity $H_{\rm C}$ were measured with a vibration sample magnetometer in magnetic fields up to 10 kOe. The constant K_4 in the (001) plane was determined by a torque magnetometer with a sensitivity of $3.76 \cdot 10^{-9}$ Nm in a magnetic field H = 0-12 kOe. The saturation magnetization $M_{\rm S}^0$ and the magnetic fourfold anisotropy constants K_4^0 were determined for the total volume (*V*) of the 32Rh/68Fe(001) and 24Rh/76Fe(001) bilayers. All saturation magnetization and the magnetic anisotropy constants measurements were carried out on the samples at room temperature using the torque method presented in Refs. [20,21].

The reversible antiferromagnetic to ferromagnetic (AFM-FM) phase transitions were checked using a Magnetic Properties Measurement System, Quantum Design (MPMS-XL). Magnetic fields of H = 1.0 kOe were applied along the in-plane [100] MgO direction which coincides with the easy axis of the Rh/Fe(001) bilayers at all measurements in the 10–450 K temperature interval.

To obtain the films of the Fe_3Rh composition stoichiometry, the single crystal Fe(001) layer of the 75 nm thickness was deposited at a

temperature of 200 °C on the Fe₆₈Rh₃₂(001) samples, obtained by annealing at 800 °C for 1 h of 32Rh/68Fe(001) bilayers grown on MgO (001), as described above. The EDS analysis showed that the obtained Fe (001)/Fe₆₈Rh₃₂(001) samples have an atomic composition of Fe₇₆Rh₂₄ (close to the Fe₃Rh composition stoichiometry).

3. Results and discussion

3.1. Structural and magnetic phase transformations in the 32Rh/ 68Fe(001) bilayer during annealing up to 800 °C

Recent studies of solid-state reaction in Rh/Fe(001 bilayers [21] have shown the sequential formation of nonferromagnetic B2 (nfm-B2) at above 100 °C and the ferromagnetic B2- α_l phase above 300 °C. The B2- α_i is a modification of the α phase having a low magnetization about 825 emu/cm³ and a reversible antiferromagnetic to ferromagnetic (AFM-FM) phase transition that exists in a narrow composition range of nearly equiatomic concentration. It has been shown that in the 45Rh/55Fe(001) bilayers [21], the Fe-rich highly B2-ordered α_h phase is formed from the equiatomic α_l compound above 400 °C by the solid-state reaction $\alpha_l + Fe \rightarrow (\sim 450 \text{ °C}) \alpha_h$. The α_h phase possessed a magnetization $\sim 1270 \text{ emu/cm}^3$ and the reversible AFM-FM transition was completely suppressed.

In the experiments, we used Fe-rich 32Rh/68Fe(001) bilayers that were grown on a MgO(001) substrate with the atomic composition of Fe₆₈Rh₃₂. Fig. 1a shows a schematic diagram of the phase transformations consistently occurring in 32Rh/68Fe(001) bilayers on a MgO (001) substrate during annealing from room temperature to 800 °C, which builds on the X-ray diffraction (XRD) analysis (Fig. 1b) and magnetic measurements (Fig. 1c and d). The transformations consist of the successive formation of a nonferromagnetic B2-FeRh phase (nfm-B2) at above 100 °C in a thin layer at the Rh/Fe interface of the asdeposited sample, which above 300 °C turns into the B2- α_1 phase. Above 500 °C the synthesized Fe₆₈Rh₃₂ films contain the highly ordered B2- α_h (further down in the text B2) and new ordered D0₃ phases. The lattice constant of the ordered B2-FeRh phase (space group Pm-3m, lattice constant a = 0.2993 nm) is very close to the database data (PDF 4+ card # 01-073-2618, Pm-3m, *a* = 0.2987 nm). In the case of the ordered D0₃-Fe₃Rh(001) phase (space group Fm-3m, lattice constant a = 0.5912 nm) the crystallographic parameters were experimentally determined for the first time. As shown in Fig. 1b the nfm-B2 and ferromagnetic α_l phases have very low intensity of the (001) and (002) diffraction reflections, which indicates that both phases epitaxially grow on the Fe(001) surface and both have a chemically ordered B2 nanocrystalline structure. After annealing above 400 °C the intensity of the B2(001) and B2(002) diffraction reflections starts to grow and becomes very strong above 500 °C, which indicates the formation of high structural quality epitaxial B2(001) (B2- α_h) layers on the MgO(001) surface. The phase sequence 32Rh/68Fe \rightarrow (~100 °C) nfm-B2 \rightarrow (300 °C) α_1 \rightarrow (~450 °C) $\alpha_{\rm h}$ coincides with the phase sequence in 45Rh/55Fe bilayers [21]. However, above 500 °C, in addition to the superstructural (001) and fundamental (002) reflections from B2 (B2- α_h), new reflections appear with interplanar spacings $d_{(002)} = 0.2948$ nm and $d_{(004)} =$ 0.1474 nm, reflections that belong to a new ordered phase in the Fe-rich region of the Fe-Rh phase diagram. Below, we will show that the new phase has a Fe₃Rh composition and a DO₃ ordered Fe₃Al structure. This means that, unlike 45Rh/55Fe(001), which above 500 °C contains only the B2 phase [21], the Fe-rich 32Rh/68Fe(001) bilayers contain a mixture of ordered (B2 + D0₃) phases. The intensity of D0₃ reflections is well below the intensity from the B2 phase reflections (Fig. 1b), which implies that B2 is the dominant phase after annealing in the range from 600 $^\circ\text{C}$ to 800 $^\circ\text{C}.$ Fig. 1c shows the dependence of the in-plane relative magnetic anisotropy constant $K_4(T)/K_4^0$ and the relative magnetization $M_{\rm S}(T)/M_{\rm S}^0$ as a function of the annealing temperature (T) of the 32Rh/68Fe(001) bilayers, which are similar to the 45Rh/55Fe(001) bilayers [21]. The saturation magnetization $M_{\rm S}^0 = 1100 \text{ emu/cm}^3$ and



Fig. 1. The synthesis of ordered $B2 + D0_3$ phases in the 32Rh/68Fe(001) bilayer after annealing up to 800 °C. (a) Schematic of the phase transformations in the 32Rh/68Fe(001) bilayer showing the successive formation of the thin non-magnetic nanocrystalline B2-FeRh (nfm-B2) film, which evolves into the lowmagnetization modification α'_{l} above 300 °C. After annealing above 400 °C the α_1 phase develops into a mixture of $\alpha'_{h}(B2)$ and D0₃ phases with high magnetization. (b) XRD patterns of the 32Rh/68Fe(001) bilayer after annealing at different temperatures, indicating the phase transition from the nanocrystalline α_1 phase into a mixture of α_h (B2) and D0₃ phases with high magnetization above 400 °C. (c) Evolution of the magnetization $M_{\rm S}(T)/M_{\rm S}^0$ and anisotropy $K_4(T)/M_{\rm S}^0$ K_4^0 constants as a function of the annealing temperature. (d) Temperature dependent magnetization curve in 32Rh/68Fe(001) bilayers after annealing at 400 °C and 800 °C. (e) Magnetic hysteresis loop of the $Fe_{32}Rh_{68}$ film containing a mixture of the $B2 + D0_3$ phases measured after annealing at 800 °C ($M_s^0 \sim$ 1590 emu/cm³, $H_c \sim 200$ Oe, $M_r^0 \sim 600$ emu/cm³).

magnetic fourfold anisotropy constant $K_4^0 = 3.1 \cdot 10^5 \text{ erg/cm}^3$ were determined for the total volume of the 32Rh/68Fe sample. In the temperature range of 500-800 °C, the synthesized Fe₆₈Rh₃₂ reaction products (B2 + D0₃) have a saturation magnetization of \sim 1590 emu/cm³ $(M_{\rm S}(T)/M_{\rm S}^0 \sim 1.45)$, which is significantly higher than ~ 1275 emu/cm³ of the highly ordered B2 in 45Rh/55Fe(001) samples [21]. This suggests that the magnetization of the D0₃ is higher than the magnetization of the B2 phase. The $K_4(T)/K_4^0$ dependence shows that after annealing above 800 °C the magnetic fourfold anisotropy constant $K_4(B2 + D0_3) \sim 0$. The $K_4(B2 + D0_3)$ constant contains contributions from the first order magnetocrystallographic anisotropy constants $K_1(B2)$ and $K_1(D0_3)$. The value of the $K_1(B2)$ constant of the 45Rh/55Fe(001) samples is close to zero [21], thereby the anisotropy $K_1(D0_3)$ constant should be expected to have a low value. Fig. 1d shows the temperature dependent magnetization curve in 32Rh/68Fe(001) bilayers after annealing at 400 °C and 800 °C. The samples show hysteresis of the magnetization in the antiferromagnetic to ferromagnetic (AFM-FM) phase transition, which indicates the existence of a B2- α_l ' layer in 32Rh/B2- α_l '/68Fe(001) films after annealing at 400 °C [21]. However, in the temperature range of 600–800 °C, a mixture of the two (B2 + D0₃) phases occurs simultaneously, in which the AFM-FM phase transition is completely

suppressed. The magnetic (*M*-*H*) hysteresis loop after annealing at 800 °C is presented in Fig. 1e, which confirms that the (B2 + D0₃) mixture is a soft magnetic material (saturation magnetization $M_s^0 \sim$ 1590 emu/cm³, coercivity $H_c \sim$ 200 Oe, remanent magnetization $M_r^0 \sim$ 600 emu/cm³). To summarize, we show the phase evolution in the 32Rh/68Fe(001) bilayer, which like the 45Rh/55Fe(001) bilayer, contains the phase sequence (~100 °C) nfm-B2 \rightarrow (~300 °C) B2- α_l \rightarrow (~450 °C) B2- α_h , and their only difference is that above 500 °C a new ordered D0₃ phase appears in the 32Rh/68Fe films.

3.2. Structural and magnetic properties of the ordered DO₃ -Fe₃Rh phase

In the experiments, we used Fe(001)/Fe₆₈Rh₃₂(001) bilayers that were grown on a MgO(001) substrate. To obtain the films of the Fe₃Rh composition stoichiometry, a single crystal Fe(001) layer of the 75 nm thickness was deposited on the Fe₆₈Rh₃₂(001) samples at a temperature of 200 °C. The EDS analysis showed that the obtained Fe(001)/Fe₆₈Rh₃₂(001) samples have an atomic composition of Fe₇₆Rh₂₄ (close to the Fe₃Rh composition stoichiometry). The process of obtaining the Fe(001)/Fe₆₈Rh₃₂(001) bilayers is described at the end of Section 2.

Fig. 2a shows a schematic diagram of the phase transformations



Fig. 2. Formation of the D0₃-Fe₃Rh(001) basic phase in the Fe(001)/Fe₆₈Rh₃₂(001) bilayer under annealing up to 800 °C. (a) Schematic of the phase transformations in the Fe(001)/Fe₆₈Rh₃₂(001) bilayer, showing synthesis of the (B2 + D0₃) mixture above 400 °C. (b) XRD patterns of the Fe(001)/ Fe₆₈Rh₃₂(001) bilayer after annealing at different temperatures, indicating that the final reaction products contain a residual B2 phase in addition to the main D0₃ phase. (c) Temperature-dependent magnetization $M_{\rm s}(T)/M_{\rm s}^0$ and anisotropy $K_4(T)/K_4^0$ constants confirming the synthesis of D0₃-Fe₃Rh above 400 °C. (d) Magnetic hysteresis loop of the Fe₂₄Rh₇₆ film measured after annealing at 800 °C ($M_{\rm s}^0$ ~ 2150 emu/cm³, $H_{\rm c} \sim 150$ Oe, $M_{\rm f}^{\rm c} \sim 540$ emu/cm³.

consistently occurring in Fe(001)/Fe₆₈Rh₃₂(001) bilayers on a MgO (001) substrate during annealing from 400 °C to 800 °C. The main result is that the basic ordered D03-Fe3Rh(001) phase is formed in the reaction products above 400 °C, which is epitaxially intergrown with the residual B2(001) phase. Fig. 2b shows XRD patterns of the annealed samples. The original Fe(001)/Fe₆₈Rh₃₂(001) bilayers contained only reflections from the Fe(001) and Fe₆₈Rh₃₂(001) layers after annealing at 400 °C. This means that the Fe(001) grew epitaxially in a cube on cube orientation on the mixture of the ordered $B2(001) + D0_3(001)$ phases. After annealing at 400 °C, the intensity of the Fe(001) peak decreases and disappears completely above 500 °C. This means that the reaction between Fe(001) and the Fe₆₈Rh₃₂(001) films ends. In the temperature range 500–800 °C, in contrast to the 32Rh/68Fe(001) samples in which the main B2 phase is formed, the X-ray of the diffraction patterns of the 24Rh/76Fe(001) samples show $d_{(002)} = 0.29467$ nm and $d_{(004)} = 0.14734$ nm reflections from the new ordered D0₃ phase, which proves that it is the main phase and has Fe₃Rh composition. However, the B2 reflections are also present, and they do not disappear with prolonged annealing at 800 °C. This suggests that the final reaction products contain a residual B2 phase in addition to the main DO₃ phase. A similar formation of intergrown DO₃ and B2 phases was observed in FeGa films [22]. Since the D0₃ vol fraction in the Fe76Rh24 is much larger than in the Fe68Rh32 samples, it follows that the D03 structure formation is the result of the solid-state reaction (1) between Fe and B2 which starts above the initiation temperature $T_{\rm in} \sim 400^{\circ} \rm C$.

$$Fe + B2 (\sim 400^{\circ}C) \rightarrow D0_3$$
 (1)

Fig. 2c shows the dependence of the in-plane relative magnetic anisotropy constant $K_4(T)/K_4^0$ and the relative magnetization $M_S(T)/M_S^0$ as a function of the annealing temperature (T) of the Fe(001)/Fe₆₈Rh₃₂(001) films. Both dependencies begin to change after annealing at 400 $^{\circ}$ C, which is associated with the start of the reaction (1). The values $M_{\rm S}^0 = 1650 \text{ emu/cm}^3$ and $K_4^0 = 0.5 \cdot 10^5 \text{ erg/cm}^3$ for the Fe(001)/ Fe₆₈Rh₃₂(001) films contain contributions from the magnetizations and anisotropy constants of the Fe(001) and Fe₆₈Rh₃₂(001) layers, respectively. A rough estimation of the volume fraction of the B2 phase in the (B2+D0₃) mixture in the Fe₇₆Rh₂₄ samples, obtained by calculating the total intensity of the D0₃(004) and B2(002) fundamental reflections, does not exceed 10%. Synthesized Fe76Rh24(001) samples contain the basic D03 phase after annealing, therefore, neglecting the contribution from the B2 phase, the saturation magnetization $M_{\rm S}({\rm D0}_3\text{-}{\rm Fe}_3{\rm Rh})$ is 2150 emu/cm³ and the first magnetic anisotropy constant $K_1(D0_3$ -Fe₃Rh) ~ + $0.5 \cdot 10^5$ erg/cm³. The constant $K_1(D0_3$ -Fe₃Rh) was assigned the "+" sign, since the easy axis, as in single-crystal bcc-Fe, coincides with the [100] and [010] directions of the D0₃ phase. The $K_1(D0_3$ -Fe₃Rh) value has a low value, therefore it does not contribute to the magnetic fourfold anisotropy constant K_4 of the synthesized Fe₆₈Rh₃₂ films. The resulting value $M_{\rm S}$ (D0₃-Fe₃Rh) is slightly higher than the values of the Fe₇₅Rh₂₅ bulk samples ($\sim 2000 \text{ emu/cm}^3$), which were obtained in Ref. [13]. Both observations imply that rhodium atoms have a strong effect on increasing the magnetization of the D0₃-Fe₃Rh phase. Fig. 2d shows the magnetic (M-H) hysteresis loop of the Fe76Rh24 film after annealing 800 °C which does not saturate at magnetic fields of 2 kOe (saturation magnetization $M_{\rm s}^0 \sim 2150 \text{ emu/cm}^3$, coercivity $H_{\rm c} \sim 150$ Oe, remanent magnetization $M_{\rm r}^0 \sim 540 \text{ emu/cm}^3$). One of the possible reasons for the increase of magnetization at higher fields is a noncollinear alignment of the Fe magnetic moments in the D0₃ phase. Thus, soft magnetic D0₃-Fe₃Rh with a saturation magnetization of 2150 emu/cm³ and a magnetic anisotropy constant $K_1 \sim +0.5 \cdot 10^5 \text{ erg/cm}^3$ is the basic phase in Fe₇₆Rh₂₄ films.

3.3. Epitaxial Growth of D0₃-Fe₃Rh(001) on MgO(001)

The results of the asymmetrical XRD φ -scan for Fe₆₈Rh₃₂ and Fe76Rh24, containing reflections after annealing at 800 °C are presented in Fig. 3a,b. As is known, many Fe-rich alloys with cubic crystal lattices of 3:1 stoichiometry, such as Fe-Al, Fe-Ga, Fe-Si, Fe-Pd, Fe-Pt and etc. can crystallize into L1₂ or D0₃ ordered structures. In Fig. 3c, the φ -2 θ scans show that only the (00l) reflections of B2 and the new phase are visible, along with the (00l) reflections of MgO, which indicates that the (001) planes of B2 and the new phase are parallel to the (001) planes of MgO. Previously was shown, using φ -2 θ asymmetrical diffraction images, that B2(001)[001] || MgO(001)[110] is the epitaxial orientation for the B2 phases grown by the solid-state reaction between the Rh and Fe(001) layers on the MgO(001) substrate [21]. Asymmetrical φ -scans of the (311) reflection for $L1_2$ (expected lattice constant a = 0.367 nm [13]), and the (112) reflection for D0₃ were used to find the orientation relationships and ordering types in the new phases $(D0_3 \text{ or } L1_2)$ (Fig. 3c). The φ -scans were carried out using the (311) reflection having no fixed points, which indicates the absence of $L1_2$ ordering in the new phase. Because the new phase has Fe₃Rh composition the φ -scans of the (112) reflection determined that the new phase must have an ordered D03 crystal structure (space group Fm-3m). Asymmetrical diffraction images on Fig. 3c define the orientation relationship between the D03-Fe3Rh film and the MgO substrate, schematically shown in Fig. 3d: D0₃(001) [001] || MgO(001)[110]. Lattice parameters were determined for the $Fe_{68}Rh_{32}$ ($\alpha_{B2} = 0.2993$ nm, $\alpha_{D03} = 0.5912$ nm) and $Fe_{76}Rh_{24}$ ($\alpha_{B2} =$ 0.2980 nm, $\alpha_{D03} = 0.5888$ nm) samples after annealing at 800 °C.

3.4. Ordered D0₃-Fe₃Rh phase and the Fe–Rh phase diagram

The Fe-rich part of the Fe–Rh phase diagram [16] at low temperatures <600 °C contains 0–20 at.% of Rh in the chemically disordered bcc (A2) alloys and between 20 and 50 at.% in the chemically ordered bcc (B2) alloys. At high temperatures above 600 °C only the chemically disordered fcc (A1) alloys and paramagnetic B2 exist in the concentration range of 0–50 at.% Rh. Fig. 4 shows the modified Fe-rich part of the Fe–Rh phase diagram based on the experimental results of this work. The region of the phase mixture of (D0₃ + B2) is limited by the compositions of Fe₇₆Rh₂₄ and Fe₆₈Rh₃₂. However, the real (D0₃ + B2) region can be much wider, since the phase boundaries A2/(D0₃ + B2) and (D0₃ + B2)/B2 have not been strictly found. The characteristic temperature 400 °C was also detected for the (D0₃ + B2) region by magnetic measurements (Figs. 1c and 2c).

For many film systems, it has been shown that the initiation temperature T_{in} of solid-state reaction in A/B bilayers starts at the temperature of phase transformations T_k in A-B binary system (see for example [21–27] and the references therein for more details about this topic). Consequently, the equality $T_{in} = T_k$ can be used to analyze the phase equilibrium diagrams. This implies that the study of reaction in A/B bilayers with different layer ratios is the study of the A-B phase diagram as the temperature increases from room temperature to high temperatures (bottom-up approach). This approach differs from the generally



Fig. 4. Modified phase diagram of an Fe–Rh system adapted from ref. 16. Revised Fe-rich part of the Fe–Rh phase diagram contains a new $B2+D0_3$ region. The real $B2+D0_3$ region can be wider, since the phase boundaries $A2/B2+D0_3$ and $B2/B2+D0_3$ are not precisely defined.





accepted construction of the equilibrium phase diagram of binary metallic systems by cooling the melt to room temperature (top-down approach). Therefore, the temperature hysteresis of the phase formation is possible at a given concentration via both approaches. As argued above, it is assumed that the initiation temperature of solid-state reaction (1) $T_{\rm in} = 400$ °C coincides with the temperature of the unknown phase transformation in the Fe–Rh system. The bulk phase diagram for Fe–Rh above 600 °C contains an A1 phase with an fcc structure, however, A1 is not observed in the Fe₇₆Rh₂₄ and Fe₆₈Rh₃₂ compositions up to 800 °C. We believe that the bottom-up approach suggests the formation of the A1 phase occurs at temperatures above 800 °C. Thus, additional studies are needed to refine the phase boundaries of the D0₃ + B2 region and the high-temperature part of the Fe-rich part of the Fe–Rh phase diagram.

4. Conclusions

In summary, we have discovered a new ordered D03-Fe3Rh phase, using a solid-state reaction between epitaxial Fe(001) and polycrystalline Rh layers. The epitaxially intergrown mixture of the $DO_3(001) + B2(001)$ phases grows epitaxially on the MgO(001) substrate in both 32Rh/68Fe(001) and 24Rh/76Fe(001) bilayers above 400 °C. The high structural perfection of the ordered D0₃(001) and B2 (001) phases obtained by this method makes it possible to separate close reflections from the D03-Fe3Rh and B2-FeRh phases. The lattice parameter $\alpha_{D03} = 0.5888$ nm, the saturation magnetization $M_{\rm S}({\rm D0}_{3}$ - $Fe_3Rh) \sim 2150 \text{ emu/cm}^3$ and the first magnetic anisotropy constant $K_1(D0_3$ - Fe₃Rh) ~ + 0.5 · 10⁵ erg/cm³ were determined for the D0₃-Fe₃Rh phase. Thus, our work not only justified a method for finding the ordered D0₃-Fe₃Rh (001) phase, but also suggests that this method can be extended for iron-based alloys to find ordered D03 structures in which D03 has closely spaced X-ray diffraction reflections with A2 and B2 phases.

CRediT authorship contribution statement

All authors contributed to the manuscript and the interpretation of the data.

Author statement

All authors acknowledge that the material presented in this manuscript has not been previously published, nor is it simultaneously under consideration by any other journal.

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

Data availability

No data was used for the research described in the article.

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